Near-Infrared C$_2$H$_2$ Detection System Based on Single Optical Path Time Division Multiplexing Differential Modulation Technique and Multi-Reflection Chamber

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Featured Application: Near-infrared ppm-level C$_2$H$_2$ detection with time division multiplexing differential modulation, suitable for the precision improvement of other near-infrared single optical path gas detection systems.

Abstract: A time division multiplexing differential modulation technique is proposed to address the interference problem caused by the fluctuation of laser light intensity in the single optical path detection system. Simultaneously, a multi-reflection chamber is designed and manufactured to further improve the system’s precision with an optical path length of 80 m. A near-infrared C$_2$H$_2$ detection system was developed. The absorption peak of the acetylene (C$_2$H$_2$) molecule near 1520 nm was selected as the absorption line. A laser driver is developed, and a lock-in amplifier is used to extract the second harmonic (2f) signal. A good linear relationship existed between C$_2$H$_2$ concentration and the 2f signal, and the correlation coefficient was 0.9997. In the detection range of 10–100 ppmv, the minimum detection limit was 0.3 ppmv, and the precision was 2%. At 50 ppmv, C$_2$H$_2$ and continuous detection for 10 h, the data average was 50.03 ppmv, and the fluctuation was less than ±1.2%. The Allan variance method was adopted to evaluate the long-term characteristic of the system. At 1 s of integration time, the Allan deviation was 0.3 ppmv. When the integration time reached 362 s, the Allan deviation was 0.0018 ppmv, which indicates the good stability of the detection system.

Keywords: near-infrared; C$_2$H$_2$ detection; TDLAS; time division multiplexing differential modulation; a multi-reflection chamber

1. Introduction

Acetylene (C$_2$H$_2$) is one of the most basic raw materials for industrial production. C$_2$H$_2$ is burned at high temperature to weld metals. C$_2$H$_2$ is easily decomposed, burned, and exploded, because its chemical character is very active. The explosion limit of C$_2$H$_2$ in air is between 2.3% and 72.3%. C$_2$H$_2$ has a lower explosion limit and a wider explosion range than other flammable and explosive gases. Therefore, highly sensitive and real-time detection of C$_2$H$_2$ concentration is especially necessary [1–4].

In recent years, tunable diode laser absorption spectrometry (TDLAS) has been widely used as a real-time and efficient online detection technique [5–9]. C$_2$H$_2$ detection is mostly performed in the near-infrared band because the laser process in the near-infrared band is mature. The reported
TDLAS-based C\textsubscript{2}H\textsubscript{2} detection system uses a wavelength modulation spectroscopy (WMS) technique to improve the signal-to-noise ratio (SNR) [10–12]. In 2018, Ming D et al. [13] developed a dual-optical C\textsubscript{2}H\textsubscript{2} detection system near 1533 nm by using WMS, with minimum detection limit (MDL) of 7.9 ppmv at 6 m and 4 ppmv at 20 m. In the actual detection, to further eliminate the influence of the fluctuation of laser light intensity, the dual optical path differential method is used for detection [14,15]. For example, in 2018, the C\textsubscript{2}H\textsubscript{2} detection system near 1534 nm, which was developed by He Q X et al. [15], has low MDL of 3.97 ppmv at 0.3 m and a precision of 3%. However, this dual optical path differential structure reduces the final detected optical power and increases system complexity. This study proposes a time division multiplexing differential modulation technique without changing the overall structure of the single optical path detection system. The DFB laser is driven and modulated by superimposing a high-frequency sine signal by a stair-stepping segmented low-frequency signal. Thus, it can eliminate the background noise caused by the fluctuation of laser light intensity, which further improves the detection precision. In addition, this study proposes the design and manufacture of a multi-reflection chamber that increases the absorption optical path and further improves the detection precision. This study proposes a time division multiplexing differential modulation technique without changing the overall structure of the single optical path detection system. The DFB laser is driven and modulated by superimposing a high-frequency sine signal by a stair-stepping segmented low-frequency signal. Thus, it can eliminate the background noise caused by the fluctuation of laser light intensity, which further improves the detection precision. In addition, this study proposes the design and manufacture of a multi-reflection chamber that increases the absorption optical path and further improves the detection precision. The precision of the single-optical TDLAS gas detection system can be greatly improved by adopting a time division multiplexing differential modulation technique and the multi-reflection chamber method. It can better satisfy the high-precision detection requirements for C\textsubscript{2}H\textsubscript{2} gas in practical applications without increasing the complexity of the system. The precision of the single-optical TDLAS gas detection system can be greatly improved by adopting a time division multiplexing differential modulation technique and the multi-reflection chamber method. It can better satisfy the high-precision detection requirements for C\textsubscript{2}H\textsubscript{2} gas in practical applications without increasing the complexity of the system.

A near-infrared DFB laser is used to detect concentrations of C\textsubscript{2}H\textsubscript{2} based on the TDLAS technique. The emitting peak wavelength of the laser is 1520 nm. A laser driver, which uses a time division multiplexing differential modulation technique to eliminate interference caused by the fluctuation of laser light intensity in the single optical path detection system, is developed. A laser temperature controller based on ARM was developed with a precision of 0.001 °C. The 2f signal's amplitude (SA) is extracted by the lock-in amplifier, and the experimental data can be processed by the algorithm to obtain the detection result. Gas detection experiments are carried out to study the response performance of the system. The system measures different concentrations of C\textsubscript{2}H\textsubscript{2} with good linearity, and the MDL is 0.3 ppmv. In addition, the detection precision and stability of the system were analyzed with a precision of 2% and ranged from 10 to 100 ppmv, and the detection data fluctuation was less than ±1.2%. Conclusions were drawn based on the experimental results. The stair-stepping segmented a low-frequency signal that can be used to achieve high-precision C\textsubscript{2}H\textsubscript{2} detection in a single optical path system, and there is no requirement for phase consistency compared with the 2f/1f normalization [16,17].

2. TDLAS-WMS Detection Theory

2.1. Absorption Line of C\textsubscript{2}H\textsubscript{2}

The transition between the ro-vibrational energy levels of a molecule need to absorb the energy of the infrared region, and it will form an infrared spectrum [18]. From the point of view of quantum mechanics, the transition between molecular energy levels satisfies the quantization conditions. Since the energy of the ro-vibrational level is discontinuous, when a level transition occurs, a specific molecule will only absorb a specific wavelength infrared light, which ensures the selectivity of the infrared absorption spectrum gas detection technique. The absorption peak of the C\textsubscript{2}H\textsubscript{2} molecule in the mid-infrared band is evidently stronger than in the near-infrared band [19]. However, the process of the near-infrared laser is mature, widely used, and relatively inexpensive. In the actual detection process, the optical path coupling of the mid-infrared laser is more difficult. Therefore, this study selects the characteristic absorption line at the near-infrared band of the C\textsubscript{2}H\textsubscript{2} molecule to detect its concentration.

The absorption spectrum of the C\textsubscript{2}H\textsubscript{2} molecule in the range of 1520.05 to 1520.15 nm is shown in Figure 1, according to the high-resolution transmission molecular absorption database 2016 (HITRAN 2016). Evidently, the line intensity of the C\textsubscript{2}H\textsubscript{2} molecule is the largest near 1520 nm at 10^{-20}. The line
intensity of common gas molecules, such as CO₂, H₂O, and CH₄, is less than 10⁻²³ at 1520 nm, which is far less than that of the C₂H₂ molecule. The interference generated can be ignored in the actual detection. Therefore, these disturbances can be ignored in the actual detection. The absorption line at 1520 nm of the C₂H₂ molecule is used comprehensively. The line intensity is set to 1.340 × 10⁻²⁰ to obtain accurate detection results and avoid interference of other gas molecules as much as possible, considering these factors.

![Absorption line selected at 1520.1 nm](image)

**Figure 1.** C₂H₂ (blue) absorption line in the range of 1520.05–1520.15 nm. The absorption line strength of CO₂, H₂O, and CH₄ is much smaller than that of C₂H₂.

### 2.2. Direct Absorption Spectroscopy

The Beer–Lambert law indicates that, when a beam of infrared light is transmitted through the C₂H₂ molecule, a specific wavelength of light is absorbed by the C₂H₂ molecule. The selective absorption of C₂H₂ molecules causes the attenuation of light energy, and the energy of light attenuation is proportional to the number of molecules of C₂H₂ [20,21]. The light intensity after absorption is obtained as follows.

\[
I_t = I_0 e^{-\alpha(v)CL}
\]

where \(I_0\) is the initial emitting light intensity, \(L\) is the optical path length, \(C\) is the measured gas concentration, and \(\alpha(v)\) is the absorption line shape function. At a standard atmospheric pressure, we use the Lorentz line shape function to describe the gas absorption coefficient, as follows.

\[
\alpha(v) = \frac{\alpha_0}{(\frac{v-v_{cm}}{\Delta v})^2 + 1}
\]

where \(\alpha_0\) is the absorption coefficient at the center of the gas absorption line, \(v_{cm}\) is the center frequency of the gas absorption peak, and \(\Delta v\) is the half-width at half-peak of the absorption line.

### 2.3. Single Optical Path Time Division Multiplexing Differential Modulation

We can tune the laser output wavelength by modulating the drive current based on WMS [22]. A stair-stepping segmented low-frequency signal (10 Hz) was used to modulate the laser’s injection current to further eliminate the fluctuation of laser light intensity on the single-light path detection system. As shown in Figure 2, the stair-stepping segmented low-frequency signal (10 Hz) is divided into L, M, and H segments. The M segment is a sawtooth signal. This segment is used for C₂H₂ concentration detection. The laser injection current is modulated by superimposing a high-frequency sine signal (5 kHz). Thus, the laser’s output wavelength sweeps the absorption peak of C₂H₂. The currents of L and H segments are constant, and the laser’s output wavelength driven by the currents is in the non-absorption region of the C₂H₂ molecule. The laser’s mean value distribution can be considered in the short single period detection time of 100 ms at 10 Hz because its intensity fluctuation is mainly derived from its own background noise. Therefore, the average light intensity corresponding to the L and H segment waveforms is obtained, and the data processing is performed.
as the inversion basis of the light intensity \( I_{0\text{M}} \) of the M segment signal, which further eliminates the interference caused by the laser light intensity fluctuation in the single optical path detection system.

After modulation, the output frequency \( v \) and the output intensity \( I_0' \) are obtained as follows.

\[
v = v_0 + v_m \cos \omega t \tag{3}
\]

\[
I_0' = I_0(1 + \eta \cos \omega t) \tag{4}
\]

where \( v_0 \) is the center frequency of the laser’s output, \( v_m \) is the modulated SA, \( \eta \) is the light intensity modulation index, and \( f = \omega / 2\pi \) is the modulation frequency.

According to Equation (4) and Equation (1), the light intensity with gas absorption is the following.

\[
I(v, t) = I_0(1 + \eta \cos \omega t)e^{[-\alpha(v)CL]} \tag{5}
\]

The actual gas absorption and light intensity modulation are extremely small, i.e., \( \alpha(v) \ll 1 \) and \( \eta \ll 1 \). The Taylor series expansion is performed, according to Equation (5). The light intensity is obtained, ignoring the high-order terms, as follows.

\[
I(v, t) = I_0[1 + \eta \cos \omega t - \alpha(v_0 + v_m \cos \omega t)]CL \tag{6}
\]

From Equation (2), the center frequency \( v_0 \) of the laser is adjusted to coincide with the gas absorption peak, i.e., \( v_0 = v_c \). Then, the light intensity with gas absorption is obtained as follows.

\[
I(v, t) = I_0(1 + \eta \cos \omega t - \frac{\alpha_0 CL}{1 + m^2 \cos^2 \omega t}) \tag{7}
\]

\[
m = \frac{v_m}{\Delta v} \tag{8}
\]

The detection signals \( H_L(t), H_M(t), \) and \( H_H(t) \) of the photodetector can be expressed as follows.

\[
H_L(t) = spI_{0L}(1 + \eta \cos \omega t) \tag{9}
\]

\[
H_M(t) = spI_{0M}(1 + \eta \cos \omega t - \frac{\alpha_0 CL}{1 + m^2 \cos^2 \omega t}) \tag{10}
\]

\[
H_H(t) = spI_{0H}(1 + \eta \cos \omega t) \tag{11}
\]

where \( s \) is the conversion factor of the photodetector, \( p \) is the amplification factor of the detection circuit, and \( I_{0L}, I_{0M}, I_{0H} \), and are the incident light intensity corresponding to the L, M, and H segments of
the low-frequency signal that satisfy $I_{0M} = \frac{I_0 + I_{as}}{2}$. Differential signal $H(t) = \frac{H_L(t) + H_M(t)}{2} - H_M(t)$ is obtained as follows.

$$H(t) = \frac{sp I_{0M} \alpha_0 CL}{2} \left( 1 + m^2 \cos^2 \omega t \right)$$  \hspace{1cm} (12)

The Fourier series expansion is performed on Equation (12) to obtain the second harmonic component coefficient $H_{2f}$ as follows.

$$H_{2f} = \frac{ksp I_{0M} CL\alpha_0}{2}$$

$$k = \frac{2(-2 - m^2 + 2\sqrt{1 + m^2})}{m^2 \sqrt{1 + m^2}}$$  \hspace{1cm} (14)

From Equation (13), given that $m$, $L$, and $I_{0M}$ are constant, the gas concentration ($C$), and the amplitude of the $2f$ signal ($H_{2f}$) are proportional. In the experiment, the $2f$ signal is extracted, and the concentration of $C_2H_2$ can be calculated. The relationship between $H_{2f}$ and $C$ can be obtained through calibration experiments.

3. System Configuration

The $C_2H_2$ detection system is mainly divided into two components, which includes optical and electrical modules. The system structure is shown in Figure 3a. The optical module includes a DFB laser, a self-developed multi-reflection gas chamber, and a photodetector. The electrical module includes a laser driver, a high-precision temperature controller, a lock-in amplifier, and a data processing module. The design of the $C_2H_2$ detection system is shown in Figure 3b.

Figure 3. (a) Structure of the $C_2H_2$ detection system. (b) $C_2H_2$ concentration detection experimental platform.

3.1. Testing of Laser Performance

The near-infrared DFB laser (1520 nm, Nanoplus) is driven by a self-developed driver. The output wavelength of the laser is tuned by controlling the drive current and temperature, and the laser’s temperature was set to 20 °C, 25 °C, and 30 °C by the self-developed temperature controller. As the driving current increases (20 to 150 mA), the emitting peak wavelength gradually increases, and the current tuning factor is approximately 0.011 nm/mA. At a stable driving current, the emitting peak wavelength increases with temperature, and the temperature tuning factor is approximately 0.113 nm/°C. The tuning characteristics of the laser are shown in Figure 4a.

The driving current is set to 80 mA, and the temperature is set to 25 °C with a monitoring experiment for 4 h (8:30–12:30). The laser’s optical emission spectrometry is shown in Figure 4b. The emitting peak wavelength varies by approximately 0.001 nm (between 1520.011 and 1520.012 nm).
The monitoring experiment results show that the laser’s emission peak wavelength slightly changes, and the working state is stable during a long-term continuous operation.

**Figure 4.** (a) Curve of the laser’s emitting peak wavelength versus the driving current at 20 °C, 25 °C, and 30 °C. (b) Curve of the laser’s emitting peak wavelength during four hours of continuous operation.

### 3.2. Multi-Reflection Chamber

The spots of the traditional optical chamber are roughly distributed in a circular shape on the spherical mirror, as shown in Figure 5a. When the distance between the mirrors is extremely close (i.e., the physical size is small), the optical path length cannot easily satisfy the detection requirements. To improve the detection sensitivity, the proposed gas chamber uses three mirrors with a mirror diameter of 60 mm. There are two semi-circular mirrors on the front and a circular mirror on the back. The mirror surface is coated with a silver film by a vacuum coating process. The effective reflection spectrum ranges from 450 nm to 20 μm and the reflectance is greater than 95%. The optical path length can be changed by manually adjusting the incident optical path. As shown in Figure 5b, the reflected spots on the mirror surface of the gas chamber are distributed in 13 rows, which increases the utilization of the mirror surface.

**Figure 5.** (a) Mirror spots of the conventional optical chamber distributed in a circular shape. (b) Mirror spots of a multi-reflection chamber distributed in 13 rows. (c) Photo of the multi-reflection chamber.

Since the mirror is made of fused silica (JGS1) material, and the main body of the gas chamber is made of 316 steel. Therefore, the material is easy to obtain, and the cost is relatively low compared with the products of the same market. The chamber adopts a ferrule connector/asperity polishing
connector (FC/APC) fiber coupling interface, the light is coupled into the gas chamber by the fiber, and the docking is convenient. The chamber’s actual structure is shown in Figure 5c, and the optical path length is set to 80 m with 250 reflections by adjusting the incident optical path.

4. Experimental Process and Discussion

In the experiment, the laser temperature is set to 25 °C, and the emission peak wavelength is near 1520 nm. The light is coupled through a fiber into a gas chamber. We used the Gas Dilution System 4040 gas distribution system. The precision and repeatability of the instrument are ±1.0% and ±0.05%. The temperature and humidity of the laboratory are constant, and the precision of gas preparation can be controlled at ±1.2%.

4.1. Calibration Results and Response Time

To measure the relationship between the 2f SA and C2H2 concentration, 10 C2H2 samples with concentrations of 10, 20, 30, 40, 50, 60, 70, 80, 90, and 100 ppmv are prepared by dynamic gas preparation. The prepared C2H2 sample was sequentially introduced into the gas chamber at intervals of 5 min. As shown in Figure 6a, the 2f SA obtained in the experiment was expressed as max (2f). The 2f SA is linear with the C2H2 concentration, and Equation (15) is obtained by linear fitting (where C_{C2H2} is the measured C2H2 concentration).

\[
\max(2f) = 0.0241C_{C2H2} - 0.0015
\]  

Then, Equation (16) is derived as follows.

\[
C_{C2H2} = 41.4938\max(2f) + 0.0622
\]  

The 2f SA can be converted to C2H2 concentration by using Equation (16) with a linear correlation of 0.9997, and the fitted curve is shown in Figure 6b.

The response time of the system depends on the time of data processing and the speed of gas diffusion. The detection system can obtain the concentration value in one stair-stepping segmented signal cycle, and 10 concentration values are displayed after average processing. In the case of static gas preparation (i.e., using an air bag to inject C2H2 into the gas chamber), the response time is approximately 2 s.

![Figure 6](image)

**Figure 6.** (a) Relationship between 2f signal’s amplitude (SA) and concentration C2H2 (60 samples per group), and detection data recorded every 5 s. (b) Fitted curve of the 2f SA and concentration C2H2.

4.2. Detection Limit and Precision

To test the minimum detection limit (MDL), 50 ppmv C2H2 is introduced into the gas chamber. The 2f signal detected by the lock-in amplifier is shown in Figure 7a. When the C2H2 concentration is 0 ppmv, the 2f signal is as shown in Figure 7b. The 2f SA is 1.206 V when the C2H2 concentration is
50 ppmv, and the system noise’s amplitude (NA) of 0 ppmv is 0.007 V. The SNR is 44.725 dB, and the MDL is as follows: 50 ppmv × NA / SA ≈ 0.3 ppmv.

He, Q.X. et al. used a 1.534-µm tunable diode laser and a miniature gas chamber to achieve C₂H₂ detection in Reference [19], and the minimum detectable absorption is 2.66 × 10⁻³. In this paper, different concentrations of C₂H₂ were introduced into the gas chamber, and the minimum detectable concentration is 0.3 ppmv. Therefore, the minimum detectable absorption is 3.01 × 10⁻³.

To test the precision of the detection system, nine C₂H₂ samples at 10, 15, 25, 35, 45, 55, 75, 95, and 100 ppmv were prepared. The detection system calculated the concentration values according to the fitting formula. The error analysis of the detection results is shown in Figure 8. The blue line indicates the absolute error of the experimental results, and the red line reflects the relative error. The maximum absolute error is 0.3 ppmv, and the maximum relative error is within ±2%. Therefore, the precision is 2% in the detection range from 10 to 100 ppmv.

4.3. Stability Analysis of the System

To test the stability of the detection system, the C₂H₂ sample at 50 ppmv was introduced into the gas chamber, and the concentration value was recorded every second. The experiment was carried out for 10 h. The result is shown in Figure 9a. The 2f SA ranged from 1.194 to 1.218 V, and the C₂H₂ concentration was between 49.6 and 50.6 ppmv. Therefore, the fluctuation was less than ±1.2%, and the average value of the data was 50.03 ppmv.
Allan deviation was adopted to analyze the long-term characteristics of the detection system [23]. The result shown in Figure 9b indicates that the Allan deviation at the initial integration time of 1 s is 0.3 ppmv. If the integration time continues to increase, then Allan deviation is greatly reduced. Moreover, the Allan deviation can be reduced to 0.0018 ppmv when the integration time is increased to 362 s. The results indicate that the detection system has good stability.

![Figure 9. (a) Results of the stability test of 50 ppmv C₂H₂. (b) Analysis of Allan deviation for long-term characteristics of the detection system.](image)

**4.4. Comparison**

We conducted a large number of experiments using a 15-cm single-pass cell and an 80-m multi-pass cell, which clarifies the increase in detection sensitivity caused by modulation and the increase in sensitivity attributable to the multi-reflection cell. Many researchers have recently conducted in-depth research on C₂H₂ detection technology in the near-infrared band. The performance comparison between the proposed C₂H₂ detection system and the reported near-infrared C₂H₂ detection system is shown in Table 1. The C₂H₂ detection system in References [13,15,19,24] uses the conventional driving method of the TDLAS-WMS technique. The laser’s driving signal is generated by super-imposing low-frequency sawtooth waves and high-frequency sine waves. Precision and error are greater than 3%. Time division multiplexing differential modulation is used to eliminate light intensity fluctuations, and system stability has been further improved, with a lower detection limit and higher precision.

**Table 1. Performance comparison between the proposed C₂H₂ detection system and the reported C₂H₂ detection system.**

<table>
<thead>
<tr>
<th>Reference/Type</th>
<th>Detection Wavelength</th>
<th>Detection Technique</th>
<th>Detection Limit (ppmv)</th>
<th>Detection Error (%)</th>
<th>System Stability (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>[13]</td>
<td>1530 nm</td>
<td>TDLAS-WMS</td>
<td>4</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>[15]</td>
<td>1530 nm</td>
<td>TDLAS-WMS</td>
<td>3.97</td>
<td>3</td>
<td>3</td>
</tr>
<tr>
<td>[19]</td>
<td>1534 nm</td>
<td>TDLAS-WMS</td>
<td>200</td>
<td>5</td>
<td>2</td>
</tr>
<tr>
<td>[24]</td>
<td>1530 nm</td>
<td>TDLAS-WMS</td>
<td>1.46</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Single-pass un-modulation</td>
<td>1520 nm</td>
<td>DAS</td>
<td>710</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Single-pass modulation</td>
<td>1520 nm</td>
<td>TDLAS-WMS</td>
<td>142</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Multi-pass un-modulation</td>
<td>1520 nm</td>
<td>DAS</td>
<td>1.5</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Multi-pass modulation</td>
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<td>TDLAS-WMS</td>
<td>0.3</td>
<td>±2</td>
<td>±1.2</td>
</tr>
</tbody>
</table>

**5. Conclusions**

A time division multiplexing differential modulation technique is proposed to eliminate the interference caused by the fluctuation of laser source in a single optical path detection system. A new multi-reflection chamber was designed to further improve the system detection precision. A near-infrared C₂H₂ detection system was developed based on TDLAS using a 1520 nm DFB laser. A laser driver and a high-precision temperature controller were developed, and a lock-in amplifier is
used to extract the 2f signal. A number of experiments were performed using the detection system to study its performance. The experimental results indicated that the detection system has good performance indicators. The MDL is 0.3 ppmv, and the precision was 2% in the detection range of 10 to 100 ppmv and contains errors caused by dynamic gas preparation. A 10-h stability observation experiment was performed on 50 ppmv C$_2$H$_2$, and the fluctuation was less than ±1.2%. The Allan deviation value was 0.3 ppmv when the initial integration time was 1 s, according to the Allan deviation analysis experimental data. When the integration time increased to 362 s, Allan deviation value was 0.0018 ppmv, which indicated the system’s good stability. In addition, the time division multiplexing differential modulation technique and the chamber structure adopted by the system were suitable for the precision improvement of other near-infrared single optical path gas detection systems. Therefore, the proposed approach was flexible and practical.

**Author Contributions:** Conceptualization, B.W. Methodology, B.W. and H.L. (Hongfei Lu). Software, B.W. Validation, B.W., H.L. (Hongfei Lu), and C.C. Formal analysis, L.C. Investigation, H.L. (Hongfei Lu). Writing—original draft preparation, H.L. (Hongfei Lu), B.W., and C.C. Visualization, H.L. (Hongquan Lian), T.D., and Y.C. Project administration, L.C. Funding acquisition, B.W.

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**References**


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