



Article

# Facile Solution Synthesis, Processing and Characterization of n- and p-Type Binary and Ternary Bi–Sb Tellurides

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**Featured Application:** This article presents a truly scalable, energy and cost-effective, and reproducible solution chemical synthesis method for thermoelectric materials.

**Abstract:** The solution synthesis route as a scalable bottom-up synthetic method possesses significant advantages for synthesizing nanostructured bulk thermoelectric (TE) materials with improved performance. Tuning the composition of the materials directly in the solution, without needing any further processing, is important for adjusting the dominant carrier type. Here, we report a very rapid (2 min) and high yield (>8 g/batch) synthetic method using microwave-assisted heating, for the controlled growth of Bi<sub>2-x</sub>Sb<sub>x</sub>Te<sub>3</sub> (x: 0–2) nanoplatelets. Resultant materials exhibit a high crystallinity and phase purity, as characterized by XRD, and platelet morphology, as revealed by SEM. Surface chemistry of as-made materials showed a mixture of metallic and oxide phases, as evidenced by XPS. Zeta-potential analysis exhibited a systematic change of isoelectric point as a function of the material composition. As-made materials were directly sintered into pellets by using spark plasma sintering process. TE performance of Bi<sub>2-x</sub>Sb<sub>x</sub>Te<sub>3</sub> pellets were studied, where the highest ZT values of 1.04 (at 440 K) for Bi<sub>2</sub>Te<sub>3</sub> and 1.37 (at 523 K) for Sb<sub>2</sub>Te<sub>3</sub> were obtained, as n- and p-type TE materials. The presented microwave-assisted synthesis method is energy effective, a truly scalable and reproducible method, paving the way for large scale production and implementation of towards large-area TE applications.

**Keywords:** chalcogenides; microwave-assisted synthesis; polyol synthesis; thermoelectric; ZT; power factor; thermal conductivity; nanomaterial; XPS

## 1. Introduction

Metal chalcogenide-based materials are commonly good thermoelectric (TE) materials with a high thermoelectric figure of merit. Bismuth telluride ( $\text{Bi}_2\text{Te}_3$ ) and its alloys ( $(\text{BiSb})_2(\text{SeTe})_3$ ) are currently the most promising TE materials in bulk form, at ambient temperature region, due to their high figure of merit in comparison to other material [1]. The dimensionless figure of merit is defined as  $ZT = (S^2 \sigma T) / \kappa$ , where  $S$  is the Seebeck coefficient,  $\sigma$  is the electronic conductivity,  $T$  is the temperature, and  $\kappa$  is the total thermal conductivity. Theoretical analysis and experimental explorations suggest that low dimensional materials such as nanostructures, quantum wells, quantum wires, and quantum dots should have much higher thermoelectric performance than their bulk counterparts [2–4]. The highest  $ZT$  value, for  $\text{Bi}_2\text{Te}_3$  based materials, thus far ( $\sim 2.4$ ) has been obtained in  $\text{Bi}_2\text{Te}_3/\text{Sb}_2\text{Te}_3$  superlattices [5]; however, the outstanding  $ZT$  values were not reproduced, nor is the material scalable for implementing in large-scale applications. Based on the theoretical predictions, one-dimensional nanostructures including nanowires and nanotubes could reach high  $ZT$  values ( $>5$ ) [6], which has not been realized yet.

A review on the synthesis of  $\text{Bi}_2\text{Te}_3$  nanostructures has been recently reported elsewhere [7].  $\text{Bi}_2\text{Te}_3$  is normally synthesized by fusion of bismuth (Bi) and tellurium (Te) metals forming an ingot, or by powder metallurgy [8]. Arc melting, a high energy process, has been shown to be a fast and reproducible synthetic method for high performance  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  structures [9,10], even combined with solidification under magnetic field, yielding textured materials [11].  $\text{Bi}_2\text{Te}_3$  nanostructures with various morphologies such as nanoparticles [12–14], nanosheets [15,16], films [17], nanowires [18], hollow nanospheres [14], nanorods [19,20], and nanotubes [21,22] have been prepared by various routes, including solvothermal [23] or hydrothermal method, polyol method [13], electrodeposition [19,20], etc. Among these methods, solvothermal or hydrothermal reaction processes are performed in autoclaves and usually require long reaction time (in the order of many hours to days). Electrodeposition of nanowires/nanorods require custom-made templates with special pore structure, which will dictate the resultant particle morphology. Some of these techniques are either very energy intensive or time consuming, still limiting the bulk production of these materials for niche applications. Therefore, developing a more convenient and rapid way for the bulk scale synthesis of  $\text{Bi}_2\text{Te}_3$  nanostructures and its alloys is still needed. Solution chemical methods may be advantageous for the synthesis of these materials at large scale if energy effective and highly reproducible synthetic methods can be developed, which may be troubling especially when the produced volumes are large.

Microwave (MW) heating has been demonstrated to be a feasible method for the synthesis of TE nanomaterials directly in the solution phase. It has been shown to make it possible to perform reactions that have been done in several hours to days in a fraction of the time using MW-assisted heating, which is very energy effective, causing an effective and rapid volume heating, avoiding the thermal fluctuations in conventional heating, creating particles with more or less similar particle size. This type of a process is also important in the field of TEs in order to systematically investigate the influence of initial particle size and morphology in the resultant microstructure and TE transport characteristics of the materials. There are few reports on the MW-assisted method for the synthesis  $\text{Bi}_2\text{Te}_3$ . Harpeness et al. attempted to synthesize  $\text{Bi}_2\text{Te}_3$  by MW-assisted polyol method [24], where they obtained  $\text{Bi}_3\text{Te}_4$  instead of  $\text{Bi}_2\text{Te}_3$ . The synthesis of  $\text{Bi}_2\text{Te}_3$  nanorods and nanoflakes in ethylene glycol by MW heating method has been reported [25] where the reaction duration ranged from 0.5 h to 6 h. Jiang et al. [14] synthesized hollow  $\text{Bi}_2\text{Te}_3$  nanospheres at 150 °C through the MW-assisted heating method. They also reported on hollow spheres, nanosaws, and sheets [26]. Qin et al. presented a MW-assisted heating route for  $\text{Bi}_2\text{Te}_3$  nanosheets and nanotubes using ethylene glycol as the solvent with a synthesis duration from 0.5 to 1 h. They demonstrated that the morphology of the resultant structure was strongly dependent on the reaction conditions, forming only nanosheets at low temperature (180 °C), transforming to nanotubes at high temperature (195 °C) [26]. Pradhan et al. reported on the MW-assisted synthesis of spherical  $\text{Bi}_2\text{Te}_3$  nanoparticles from a homogeneous solution of tartarate complex of  $\text{Bi}^{3+}$  and  $\text{Te}^{4+}$  in a reaction time of 10 min [27].

Despite the recent developments in the synthesis of thermoelectric  $\text{Bi}_2\text{Te}_3$  nanostructures using the MW-assisted heating, only one of the above listed works reported on transport properties of the fabricated materials [27], leaving a gap from the synthesis to application. This could be mainly due to the quantity of material synthesized per batch being rather limited, as can be assessed from the experimental details (concentrations and volumes of the precursors used; see Table S1 for overview of the literature on MW-assisted synthesis of  $\text{Bi}_2\text{Te}_3$ ). It is also well-known that consolidation process, and the resultant microstructure, is critical for the final performance of the thermoelectric materials. Here, we present on the MW-assisted synthesis of  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  nanoparticles from a homogeneous solution of ethylene glycol and chloride salts of Bi and Sb, while Te is complexed with TOP/TBP, producing high purity thermoelectric nanopowders in a matter of minutes due to very effective volume heating. Besides yielding large quantity (>8 g) of the investigated materials per batch, the processing of these materials towards compacts is described, resulting in high ZT values of 1.04 for  $\text{Bi}_2\text{Te}_3$  and 1.37 for  $\text{Sb}_2\text{Te}_3$  as n- and p-type materials, respectively.

## 2. Materials and Methods

### 2.1. Synthesis and Processing of Thermoelectric Nanopowders

The following chemicals were used for the synthesis reaction:  $\text{BiCl}_3$  (Sigma Aldrich, Sweden, 98%),  $\text{SbCl}_3$  (Sigma Aldrich, Sweden, 99.95%), Te powder (Sigma Aldrich, Sweden, 99.8%), ethylene glycol (EG; Sigma Aldrich, Sweden, 99%), thioglycolic acid (TGA; Sigma Aldrich, Sweden, 98%), tri-butyl phosphine (TBP; Sigma Aldrich, Sweden, 93.5%), and tri-octyl phosphine (TOP; Sigma Aldrich, Sweden, 90%). All chemicals were used as received, without further purification.

The synthesis work was performed via MW-assisted thermolysis (flexiWAVE–Milestone, 1800 Watt; equipped with multivessel high pressure rotor) reaction inside a 100 mL Teflon vessel, in 4 parallel reaction vessels at a single run. The synthesis process starts with the preparation of precursor solutions. In the first step, Te precursor solution was prepared as a complex with TBP (or TOP) by heating a mixture of 0.96 g of Te in 6 mL TBP (or 10 mL TOP) under MW field for about 90 s at 220 °C under constant stirring, and then left to cool down. For the Bi (or Sb) precursor, 1.6 g of  $\text{BiCl}_3$  was dissolved in 20 mL ethylene glycol. The stoichiometry of ternary compositions can easily be tuned by replacing the  $\text{BiCl}_3$  precursor with equivalent moles of  $\text{SbCl}_3$ . A 1-mL aliquot of TGA as a directing agent was added to the dissolved precursors, and then sonicated to obtain a homogeneous solution. The two prepared solutions of Te-TOP complex and Bi/Sb - EG were mixed and added to a 100 mL Teflon vessel and run in the MW reactor for 2 min dwell time at 220 °C under constant stirring. The product (dark powder), which is about 2 g per vessel, was then cooled down to room temperature. The as-prepared  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  phases were separated clearly from the solvent (EG). The multivessel high pressure rotor was used for the MW synthesis work (Figure S1), where only four of the fifteen reactor positions were more than enough to prepare large quantity (>8 g) of the desired materials in a single run. Particles were easily separated from the reaction mixture by centrifuging and washing with iso-propanol and acetone several times. A typical outcome of separation process is shown in Figure S2.

### 2.2. Consolidation of Thermoelectric Nanopowders Using Spark Plasma Sintering (SPS)

The powders were consolidated using the spark plasma sintering SPS with 15 mm inner diameter graphite cylinders with top and bottom punches made as well from graphite. All powders were processed with the same parameter during the SPS process to obtain pellets with similar porosity in terms of the thermoelectric transport properties of the  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  compositions. The SPS temperature was set to 400 °C with a heating rate of 30 °C /min and 1-min dwell time. Then, 50 MPa Pressure was applied gradually during the SPS process. The holding time was set to 1 min for both applied current and pressure. Finally, the pellet was cooled from 400 °C to 250 °C, while the load was decreased from 50 MPa to 0 MPa. The pellets were then removed from the die and polished in three steps to remove the contaminated surface by the graphite punches and to obtain smooth surface with 1  $\mu\text{m}$  roughness

for further processing and analysis. Table 1 summarizes the pressing parameters and the density of the resultant pellets.

### 2.3. Structural and Morphological Characterization

X-ray Powder Diffraction (XRPD) was used to identify the crystalline phase of the nanoparticles on a PANalytical X'Pert PRO, equipped with a Copper anode (Cu-K $\alpha_1$  radiation) with a step size of 0.24° in continuous mode and a scan speed of 0.04 °/s. The crystalline phases of the Bi<sub>2-x</sub>Sb<sub>x</sub>Te<sub>3</sub> system was determined by using Rietveld refinement using High-score pro program. Scanning electron microscopy (SEM; FEI Nova 200) was used for studying the microstructure and morphology of the materials.

X-ray photoelectron spectroscopy (XPS) was used, where the XPS spectra were acquired using PREVAC 426 system configuration, equipped with SCIENTA R3000 hemispherical photoelectron spectrometer and monochromatic Al K $\alpha$  source operating at 450 W. The base pressure in the analysis chamber was better than 3 × 10<sup>-10</sup> mbar. All acquired spectra were calibrated to adventitious carbon C1s at 285 eV. The overall resolution of the spectrometer during measurements was 0.65 eV as a full width of half maximum (FWHM) of the Ag3d<sub>5/2</sub> line. After subtraction of the Shirley-type background, the core-level spectra were decomposed into main components with mixed Gaussian–Lorentzian lines (70% G + 30% L for majority of photo-peaks) by a non-linear least squares curve-fitting procedure, using CasaXPS software. XPSPeak software was also used for deconvolution, which allows for the flexible deconvolution parameter changes.

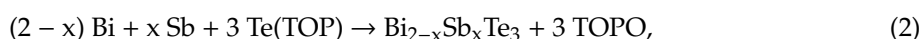
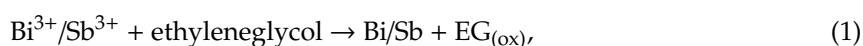
The  $\xi$ -potential of various samples were measured at room temperature with Zetasizer Nano ZS90 at an incident angle of 90°. All samples were diluted in DI H<sub>2</sub>O and titrated with 0.1 M HCl or 0.1 M NaOH to obtain pH dependent surface charge, and eventually determine the isoelectric point (*iep*). The measurements were performed in triplicate.

### 2.4. Transport Property Evaluation

The Seebeck coefficient *S* and the electrical conductivity  $\sigma$  were measured simultaneously on the pellets obtained after the SPS process using a commercial instrument ZEM-ULVAC M8 model. This system measures the *S* and  $\sigma$  based on the four-probe point method. Typical error in both electrical conductivity and Seebeck coefficient measurements is estimated as 4% in total. The total thermal conductivity  $\kappa$  was calculated using the specific heat capacity *C<sub>p</sub>*, thermal diffusivity  $\alpha$ , and density  $\rho$  using the equation  $\kappa = C_p \alpha \rho$ . The density of the samples could be obtained from Archimedes method, while a laser flash analysis system (LFA 1000, Linseis) was used for obtaining the thermal diffusivity  $\alpha$ . Differential scanning calorimetry (Netzsch-DSC 404 Pegasus214 Polyma) was used to measure the specific heat capacity *C<sub>p</sub>*. The thermal range used during the tests was 298–523 K, and with all the properties known the TE figure of merit *ZT* was estimated.

## 3. Results and Discussion

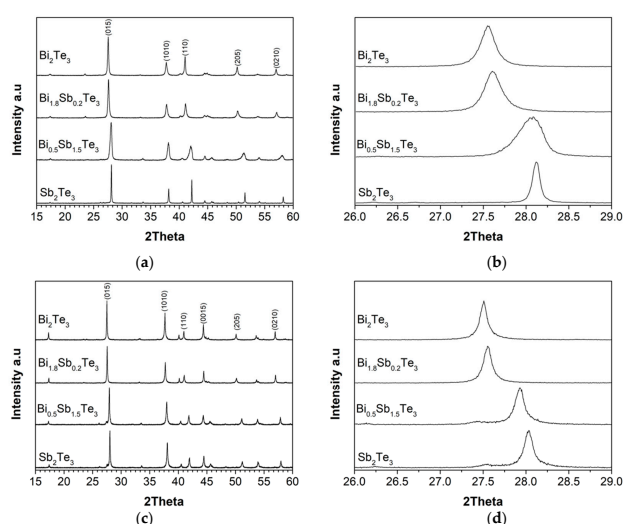
Polyol synthesis is a promising route for the synthesis of various metallic nanoparticle systems, where the solvent polyalcohol acts also as the reducing agent for the metallic ions in the solution. Here, the method esd applied for the TE material synthesis, with an important change of using MW-assisted heating for the reaction, which makes it possible to synthesize large quantities of high-quality nanomaterials in a matter of 2–3 min. The mechanism of synthesis consists of two parts, the first one being the reduction of Bi<sup>3+</sup>/Sb<sup>3+</sup> ions to metallic Bi/Sb by the solvent EG, which in turn catalyzes the breaking of the bond between Te and TOP (Te(TOP)), forming directly the crystalline Bi<sub>2-x</sub>Sb<sub>x</sub>Te<sub>3</sub> phase without needing any further thermal treatment. The overall reaction taking place during the formation of Bi<sub>2-x</sub>Sb<sub>x</sub>Te<sub>3</sub> compounds under MW irradiation can be represented as follows:



where  $EG_{(ox)}$  refers to the formation of diacetyl molecules due to the oxidation of EG during the reduction of metallic ions of  $Bi^{3+}/Sb^{3+}$  into corresponding metallic particles (see Figure S3 for molecular structure of TOP and EG as well as their oxidation products). Thermodynamic properties of binary chalcogenide compounds were reported previously by Semenkovich et al. [28], where their Gibbs free energies ( $\Delta G$ ) of formation were experimentally estimated at various temperatures. The  $\Delta G$  values are negative for pure  $Bi_2Te_3$  ( $-19.30 \pm 0.07$  kcal/mol @ 503K) and  $Sb_2Te_3$  ( $-15.4 \pm 0.3$  kcal/mol @ 540 K), revealing the feasibility and spontaneity of their formation under the reaction conditions (493 K). Chemical analysis was performed using ICP and elemental ratio, as expected from the initial precursor concentration, has been confirmed for the  $Bi_{2-x}Sb_xTe_3$  materials.

### 3.1. Structural Analysis

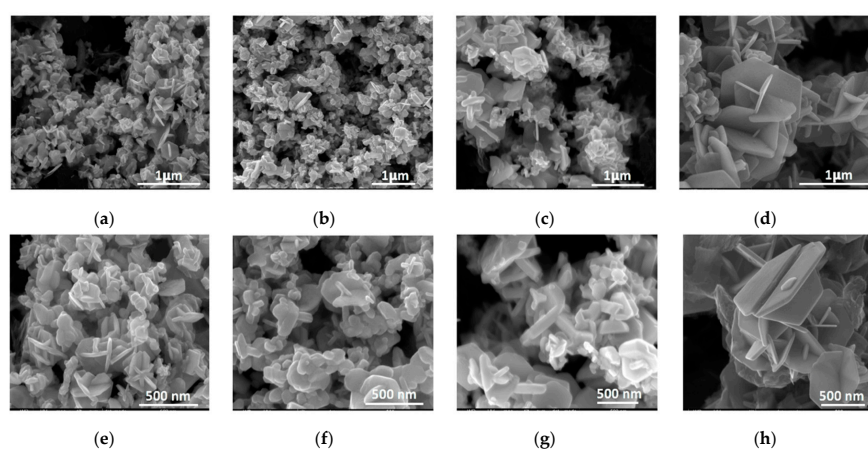
Structural analyses of as-synthesized materials and sintered pellets were performed using XRPD and the diffraction patterns are presented in Figure 1. The phases are indexed to  $Bi_2Te_3$  (ICDD: 01-082-0358) and  $Sb_2Te_3$  (ICDD: 01-071-0393) with rhombohedral crystal structure, and the corresponding indexing is designated on the powder diffraction patterns in Figure 1a. No impurity peaks are observed, revealing the high purity of the as-made materials. Due to the similarity of reduction potential of  $Bi^{3+}$  and  $Sb^{3+}$  ions, step-wise transformation of the synthetic chemistry from pure  $Bi_2Te_3$  to pure  $Sb_2Te_3$  phase was very smoothly handled at the reaction temperature. The shift of the position for the major peak indexed to (015), going from  $Bi_2Te_3$  to  $Sb_2Te_3$ , can be observed in Figure 1b, which is ascribed to systematic shrinkage of the  $Bi_2Te_3$  lattice by the introduction of Sb, as  $Sb_2Te_3$  has slightly smaller lattice dimensions ( $a = b = 4.395$  Å for  $Bi_2Te_3$  and  $4.264$  Å for  $Sb_2Te_3$ , crystal structure, as displayed in Figure S4a). Based on the lattice parameters the shrinkage in a and b corresponds to about 3% by increasing x (in  $Bi_{2-x}Sb_xTe_3$ ) from 0 to 2, while going from  $Bi_2Te_3$  to  $Sb_2Te_3$  following Vegard's rule (see Figure S4). The width of the diffraction peaks with the same index vary from sample to sample, which is ascribed to the small crystallite size and non-uniform strain within the materials. XRPD analyses performed on the sintered samples showed narrower diffraction peaks, variation of peak intensities from powder samples, due to the crystal/grain growth (Figure 1c), and reduced strain within the material upon exposure to high T and pressure for a brief period.



**Figure 1.** XRD powder patterns of powder  $Bi_{2-x}Sb_xTe_3$  samples: (a) in the  $2\theta$  range of 15–60; and (b)  $2\theta$  range of 26–29, focusing on the peak with the miller indices of (015), showing the shift of the peak position towards larger  $2\theta$  values due to shrinking lattice by substituting Bi with Sb. (c,d) The XRD patterns of the sintered pellets measured perpendicular to the pressing direction. The crystalline phases are indexed to  $Bi_2Te_3$  (ICDD: 01-082-0358) (for  $x = 0$ ) and  $Sb_2Te_3$  (ICDD: 01-071-0393) (for  $x = 2$ ) with rhombohedral crystal structure.

### 3.2. Morphology and Surface Analysis

SEM micrographs of  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  samples at different magnifications are presented in Figure 2. All samples show particles with clear hexagonal, or truncated-edge hexagonal, platelet morphology, in agreement with their rhombohedral layered crystal structure, with slight change of platelet size from one sample to another. In some cases, various platelets sharing a common axis are observed. Lateral size of hexagonal platelets varies in the range from 50 to 500 nm for different samples, where the thickness is typically around 50 nm.  $\text{Sb}_2\text{Te}_3$  sample display the largest lateral size of hexagonal platelets up to about 1  $\mu\text{m}$ , which aligns with its high crystallinity and narrower peak width observed in XRPD. A closer investigation of platelets performed by using TEM for the  $\text{Bi}_2\text{Te}_3$  sample, as an example, revealed crystalline domains (see Figure S5). There is a thin disordered layer on the nanoplatelets (see Figure S5b) which may be due to residual carbon materials or slightly different, disordered lattice.



**Figure 2.** SEM micrographs at different magnifications of as-synthesized: (a,e)  $\text{Bi}_2\text{Te}_3$ ; (b,f)  $\text{Bi}_{1.8}\text{Sb}_{0.2}\text{Te}_3$ ; (c,g)  $\text{Bi}_{0.5}\text{Sb}_{1.5}\text{Te}_3$ ; and (d,h)  $\text{Sb}_2\text{Te}_3$ .

The synthetic method and the solvents used significantly influences surface chemistry of resultant materials. In the polyol synthesis, nanoparticles may have residual chemicals on the surface, which may influence their dispersion, sintering, or some other processing characteristics. Therefore, XPS analysis was performed on the as-made materials and the resultant spectra for Bi 4f, Sb 3d, and Te 3d regions are presented in Figure S6. Carbon residue is present on all samples with varying degrees. XPS peak fitting results are summarized in Table 1, disregarding the C and O content. For all samples, during the deconvolution process, oxide phases were identified. For Sb3d and Te3d regions, metallic and oxide phases were clearly visible. In the case of Bi4f region, high resolution spectra allowed determining metallic and oxide ( $\text{Bi}_2\text{O}_3$ ) phases. The most remarkable outcome is the evidence on the presence of oxide phases on the surface of as-made nanoplatelets.

Zeta potential ( $\xi$ -potential) analysis is a technique for determining the surface charge of nanoparticles in solution in colloidal form. Nanoparticles possess a surface charge that attracts a thin layer of ions of opposite charge to their surface, and when the surface charge is balanced by the attracted oppositely charged ions, the total surface charge becomes zero. This point is defined as the isoelectric point (*iep*) and is strongly dependent on the surface exposed species.  $\xi$ -potential is usually studied as a function of an additive, most commonly as a function of pH in the solution—where pH is adjusted by the step-wise addition of acid or base solution into the colloidal system. Study of *iep* of the as-synthesized  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  materials as a function of pH is, therefore, expected to reveal the surface chemistry of as-made nanoplatelets. While all materials showed residual carbon, which should be the same type of organics due to the synthesis process used, the surface chemistry showed great variations in zeta potential analysis. Dependence of the pH at the *iep* on the Sb content,  $x$ , for  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  samples is presented in Table 1. For the pure  $\text{Bi}_2\text{Te}_3$  ( $x = 0$ ), the *iep* was reached at pH 6.3, which shifts step-wise down to 3.3 with the increasing content of Sb ( $x > 0$ ). A survey of *iep* for various oxides led to the values

of 2.9 for  $\text{Sb}_2\text{O}_3$  and 8.2 for  $\text{Bi}_2\text{O}_3$ , while no data were found for  $\text{TeO}_2$  [29]. We performed zeta potential analysis on commercial  $\text{TeO}_2$  powder and obtained an *iep* value of 5.2. XPS data, as summarized in Table 1, suggest similar  $\text{TeO}_2$  content for the samples  $\text{Bi}_2\text{Te}_3$  and  $\text{Sb}_2\text{Te}_3$ , while the other oxide phases ( $\text{Bi}_2\text{O}_3$  and  $\text{Sb}_2\text{O}_3$ ) vary in quantity depending on the *x*. The ternary phases revealed *iep* values in between those of samples where  $2 > x > 0$ . This agrees with the trend that would be expected in parallel to the changes in  $\text{Bi}_2\text{O}_3$  or  $\text{Sb}_2\text{O}_3$  content. When *x* = 0, the surface shows a closer character to  $\text{Bi}_2\text{O}_3$  with an *iep* of around 6.5, which is lower than pure  $\text{Bi}_2\text{O}_3$  due to the co-existence of  $\text{TeO}_2$ . This value changes towards more acidic region with the increasing *x* (Sb content), reaching the *iep* of 3.3 at *x* = 2, a bit higher than  $\text{Sb}_2\text{O}_3$  due to the co-presence of  $\text{TeO}_2$ . Although these data are not deterministic, they help to get a better insight into the nanoplatelets' surface in connection with the XPS results.

**Table 1.** XPS peak fitting results for Bi 4f, Sb 3d, and Te 3d regions for bulk samples, and the measured *iep* values for the  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  samples.

Samples	BE [eV] [30]	Fraction	Assigned to [31]	<i>iep</i> (pH)
$\text{Bi}_2\text{Te}_3$	157	13.11	Bi	6.3
	159	9.83	Bi–O ( $\text{Bi}_2\text{O}_3$ )	
	573	8.77	Te	
	577	19.68	Te–O ( $\text{TeO}_2$ )	
$\text{Bi}_{1.8}\text{Sb}_{0.2}\text{Te}_3$	157	6.06	Bi	6
	528.3	2.08	Sb	
	573	6.5	Te	
	159	4.56	Bi–O ( $\text{Bi}_2\text{O}_3$ )	
	530	1.42	Sb–O ( $\text{Sb}_2\text{O}_3$ )	
	577	10.35	Te–O ( $\text{TeO}_2$ )	
$\text{Bi}_{0.5}\text{Sb}_{1.5}\text{Te}_3$	157	2.38	Bi	4.5
	528.3	12.64	Sb	
	573	11.46	Te	
	159	2.38	Bi–O ( $\text{Bi}_2\text{O}_3$ )	
	530	8.43	Sb–O ( $\text{Sb}_2\text{O}_3$ )	
	577	15.7	Te–O ( $\text{TeO}_2$ )	
$\text{Sb}_2\text{Te}_3$	528.3	14.92	Sb	3.3
	573	8.38	Te	
	530	9.94	Sb–O ( $\text{Sb}_2\text{O}_3$ )	
	577	22.13	Te–O ( $\text{TeO}_2$ )	

### 3.3. Transport Property Evaluation

Nanostructured TE materials were consolidated using SPS, the critical consolidation parameters (sintering temperature, holding time, pressure, and heating rates), and characteristics of the resultant pellets are presented in detail in Table 2. All pellets possess a compaction density of about 78–83% of the theoretical density of the material, with about 20% porosity. The transport data in the temperature range 300–523 K are presented in Figure 3a–e. Table 2 also summarizes the transport data at room temperature and at the temperature where maximum *ZT* was obtained for the given material composition. No adverse effects in the transport properties were observed due to the presence of surface oxides. The Seebeck coefficient (*S*) is the key parameter to demonstrate that the materials have the transport characteristics (n- vs. p-type) planned for from the stoichiometry of the precursors. Figure 3a reveals the sign and magnitude of measured *S* values, which clearly indicates n-type character for  $\text{Bi}_2\text{Te}_3$  and  $\text{Bi}_{1.8}\text{Sb}_{0.2}\text{Te}_3$ , while p-type character was observed for  $\text{Bi}_{0.5}\text{Sb}_{1.5}\text{Te}_3$  and  $\text{Sb}_2\text{Te}_3$ . This, in agreement with literature, is a strong indication that the MW-assisted chemical synthesis process is capable of producing the phases that are designated by the initial ratio of precursor materials, without needing additional stabilizers, reflux, or other chemicals to avoid phase separation during synthesis. An interesting comparison for the obtained *S* could be for  $\text{Bi}_{0.5}\text{Sb}_{1.5}\text{Te}_3$  composition with arc-melted and textured (via solidification under magnetic field) material [11], where our *S* value is about two times

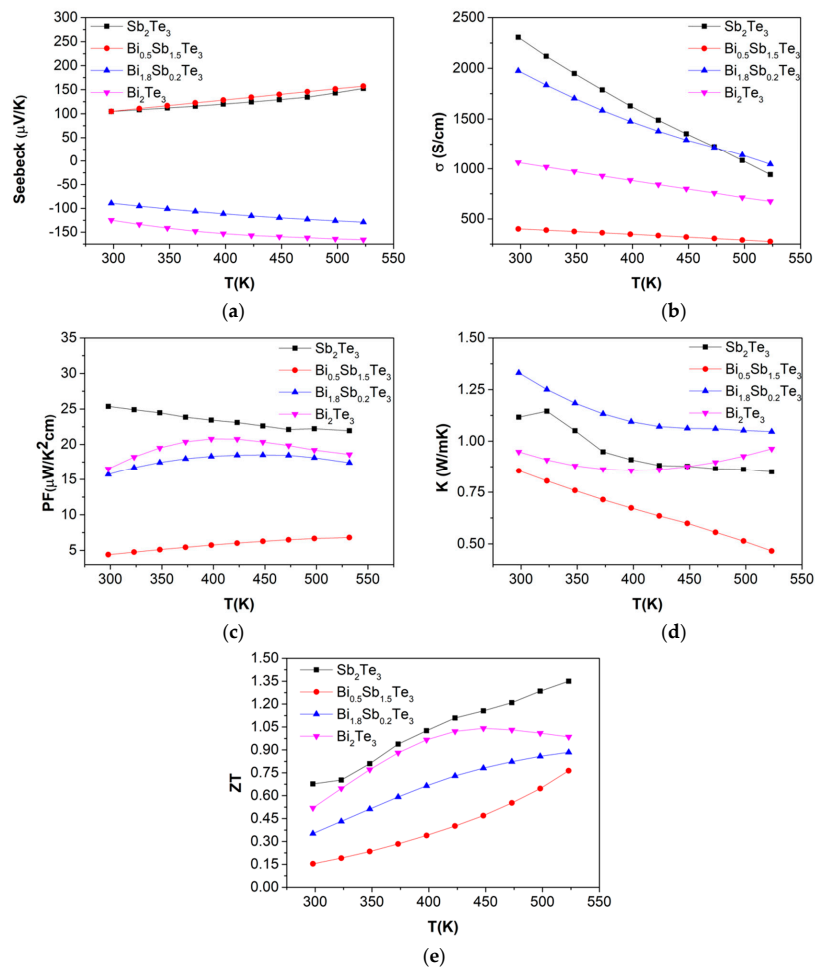
lower. Electronic transport properties are strongly correlated in classical materials, as well as by charge carrier density, defects, and scattering events. Therefore, it is possible to observe different transport values in the same composition due to the microstructure, which in turn is mostly determined by the synthesis and processing methods. The textured material possessed larger micro- to nano-sized grains, or layers, causing strong scattering of charge carriers and heat carrier phonons (decrease electronic and thermal conductivity), which in turn enhanced the *S*. SEM and TEM micrographs of our samples prepared with solution chemical process showed small nano-sized, nearly identical crystalline nanoplatelets. This microstructure allows decoupling of thermal conductivity from the electronic transport by increased density of grain boundaries, thus increased scattering of phonons (mainly reducing the lattice component of the thermal conductivity). The combined overall positive effects were seen in the details of electronic and thermal transport properties, as presented below.

**Table 2.** SPS parameters and transport properties for nanostructured bulk  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$ : sintering temperature ( $T_{\text{sint}}$ ), holding time ( $t_{\text{hold}}$ ), sintering pressure ( $P_{\text{sint}}$ ), mass density ( $\rho$ ) in percentage of the theoretical density, thermal conductivity ( $\kappa$ ), Seebeck coefficient (*S*), electrical conductivity ( $\sigma$ ), power factor (*PF*), and the thermoelectric figure of merit (*ZT*).

Sample	$T_{\text{sint}}$ (°C)	$t_{\text{hold}}$ (min)	$P_{\text{sint}}$ (MPa)	$\rho$ %	<i>T</i> (K)	$\kappa$ ( $\text{W m}^{-1} \text{K}^{-1}$ )	<i>S</i> ( $\mu\text{VK}^{-1}$ )	$\sigma$ ( $\Omega^{-1} \text{m}^{-1}$ )	<i>PF</i> ( $\mu\text{WK}^2 \text{cm}^{-1}$ )	<i>ZT</i>
$\text{Bi}_2\text{Te}_3$	400	1	50	78	300	0.94788	−124.624	106356.58	20.35	0.52
					440	0.87553	−159.466	80018.39	16.52	1.04
$\text{Bi}_{1.8}\text{Sb}_{0.2}\text{Te}_3$	400	1	50	83	300	1.33091	−89.190	197582.37	17.41	0.35
					523	1.05234	−128.942	104707.34	15.72	0.87
$\text{Bi}_{0.5}\text{Sb}_{1.5}\text{Te}_3$	400	1	50	78	300	0.85478	104.873	40077.90	6.81	0.15
					523	0.46612	157.339	27504.00	4.41	0.76
$\text{Sb}_2\text{Te}_3$	400	1	50	81	300	1.11645	104.893	230549.91	21.94	0.68
					523	0.84984	152.565	94246.95	25.37	1.37

The electrical conductivity ( $\sigma$ ) is plotted in Figure 3b, where samples show either very minor change or slight reduction of  $\sigma$  by increasing temperature, which is typical of heavily doped semiconductors, or metallic character. The highest electrical conductivity values were measured for  $\text{Sb}_2\text{Te}_3$  and  $\text{Bi}_{1.8}\text{Sb}_{0.2}\text{Te}_3$  samples. The electrical conductivity of our two n-type and one p-type  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  samples are higher than those of other solution derived samples [23,31,32] and comparable with that of p-type alloy ingot (on the level of  $1 \times 10^3$  S/cm at 300 K) [33]. The lower electrical conductivity of  $\text{Bi}_2\text{Te}_3$  in comparison to  $\text{Sb}_2\text{Te}_3$  arises from the intrinsically small band gap,  $E_g$ , of this compound ( $\sim 0.13$  eV) [34]. Given such a small  $E_g$ , the minority carriers are more easily thermally excited, and the resulting detrimental bipolar effect also introduces extra thermal conductivity. Therefore, to increase the electrical conductivity of  $\text{Bi}_2\text{Te}_3$ , the  $E_g$  has to be widened/broadened. The  $E_g$  of  $\text{Sb}_2\text{Te}_3$  is  $\sim 0.28$  eV [35], and Sb doping in  $\text{Bi}_2\text{Te}_3$  helps to optimize the carrier concentration and introduces strong point defect scattering of heat-carrying phonons. However, there is an inherent restriction of Sb doping: doping with too high a content of Sb will lead to too great a hole concentration because of the reduced formation energy of antisite defects, resulting in a reduced  $\kappa$  and an increased  $\kappa_e$  [36]. This is the main reason that  $\text{Bi}_{1.5}\text{Sb}_{0.5}\text{Te}_3$  shows both higher electrical and thermal conductivity. Similarly, Bi doping in  $\text{Sb}_2\text{Te}_3$  reduces the electrical conductivity likely due to decreased energy band gap and increased number of thermally activated minority carriers. As a result, doping with either Bi or Sb in Bi–Sb–Te compounds finely tunes the key electrical parameters such as the valley degeneracy, carrier concentration, effective mass, deformation potential, and band gap [37,38].





**Figure 3.** The temperature dependent transport properties of the synthesized  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  samples: (a) Seebeck coefficient,  $S$ ; (b) electrical conductivity,  $\sigma$ ; (c) Power factor; (d) thermal conductivity,  $\kappa$ ; and (e) the overall thermoelectric figure of merit,  $ZT$ .

The power factor (PF) ( $S^2 \sigma$ ) was calculated for all samples, as presented in Figure 3c. The highest PF values were obtained for the pure phase samples of  $\text{Bi}_2\text{Te}_3$  and  $\text{Sb}_2\text{Te}_3$ , followed by  $\text{Bi}_{1.8}\text{Sb}_{0.2}\text{Te}_3$  and  $\text{Bi}_{0.5}\text{Sb}_{1.5}\text{Te}_3$ . The magnitude of PF is comparable to or higher than that of earlier reports: the maximum power factor at 300 K we achieved are  $20.35 \mu\text{W cm}^{-1} \text{K}^2$  for  $\text{Bi}_2\text{Te}_3$  and  $21.94 \mu\text{W cm}^{-1} \text{K}^2$  for  $\text{Sb}_2\text{Te}_3$  pellet, which are much higher than other solution-derived samples—with maximum power factor  $1\text{--}9 \mu\text{W cm}^{-1} \text{K}^2$  [27,31].

Thermal conductivity ( $\kappa$ ) values were calculated from the thermal diffusivity, density, and  $C_p$  values. The trend is a decrease in  $\kappa$  by increasing temperature, while the values at room temperature (RT) range from about 0.9 to 1.2 W/mK. The sample with the composition  $\text{Bi}_2\text{Te}_3$  showed a decreasing trend of  $\kappa$ , which increased slightly after 400 K. This was attributed to bipolar effect, likely due to lower  $E_g$ , thus thermally activated minority carriers in this compound. Zhang et al. reported  $\kappa$  values in the range of 0.9–2.4 W/mK at RT for solvothermally synthesized nanoplatelets of  $\text{Bi}_x\text{Sb}_{2-x}\text{Te}_3$  [23]. Our values are similar to, or slightly lower than, their results and even that of bulk ingot  $\text{Bi}_2\text{Te}_3$  ( $1.5 \text{ W m}^{-1} \text{K}^{-1}$ ). Thermal conductivity has two contributors: electronic ( $\kappa_{\text{el}}$ ) and lattice ( $\kappa_{\text{lat}}$ ) components. Nanostructuring is known to reduce the thermal conductivity of the material by phonon scattering, which influences the  $\kappa_{\text{lat}}$  mainly.  $\kappa_{\text{el}}$  can be estimated using Wiedemann–Franz law, by calculating the electronic contribution ( $\kappa_{\text{el}}$ ) to the  $\kappa$  and thereafter  $\kappa_{\text{lat}}$  can be obtained subtracting

$\kappa_{el}$  from the total  $\kappa$  ( $\kappa_{tot}$ ). The Lorenz number ( $L_0$ ) for all compounds was calculated by using the following equations:

$$L = \left(\frac{k_B}{e}\right)^2 \left[ \frac{(r+7/2)F_{r+5/2}(\xi)}{(r+3/2)F_{r+1/2}(\xi)} - \left( \frac{(r+5/2)F_{r+3/2}(\xi)}{(r+3/2)F_{r+1/2}(\xi)} \right)^2 \right] \quad (3)$$

where  $r$  is the scattering parameter,  $k_B$  the Boltzmann constant,  $e$  is the electron charge, and  $F_n(\xi)$  is the Fermi integral given by

$$F_n(\xi) = \int_0^\infty \frac{\chi^n}{1 + e^{\chi - \xi}} d\chi \quad (4)$$

Here,  $\xi$  is the reduced Fermi energy that can be calculated from the Seebeck coefficient  $S$  and the scattering parameter  $r$  according to

$$S = \mp \frac{k_B}{e} \frac{(r+5/2)F_{r+\frac{3}{2}}(\xi)}{(r+3/2)F_{r+\frac{1}{2}}(\xi)} - \xi \quad (5)$$

We used parabolic band model and assumed the system to be highly degenerate and scattering is dominated by acoustic phonons. The calculated  $L_0$  numbers, along with the estimated  $\kappa_{el}$  and  $\kappa_{tot}$  for all samples are displayed in Table S2. Looking at the proportion of the  $\kappa_{el}$ , the major part of  $\kappa_{tot}$  is dominated by electronic contribution in  $\text{Bi}_2\text{Te}_3$  and  $\text{Bi}_{1.8}\text{Sb}_{0.2}\text{Te}_3$  samples, while  $\kappa_{lat}$  is predominant in thermal conductivity of  $\text{Sb}_2\text{Te}_3$  and  $\text{Bi}_{0.5}\text{Sb}_{1.5}\text{Te}_3$ . This might be due to either phonon-point defect scattering or phonon grain boundary scattering (due to nanosized grains). The negative value of  $\kappa_{lat}$  of  $\text{Sb}_2\text{Te}_3$  at room temperature is not realistic, which indicates that the parabolic band model may not be suitable for calculating  $L_0$  for this sample. In any case, it is important to see how  $\kappa_{lat}$  can strongly affect the transport properties and thus play a major role in enhancing the  $ZT$  value for these compositions.

Using all these transport parameters, the resultant figure-of-merit,  $ZT$ , has been estimated and the results are presented in Figure 3e, where a monotonic increase was observed for all samples by increasing temperature. Pure  $\text{Bi}_2\text{Te}_3$  and  $\text{Sb}_2\text{Te}_3$  samples show  $ZT$  values around 0.5–0.6 at RT, reaching beyond unity around 440 K. From this point onwards, the trend in  $\text{Sb}_2\text{Te}_3$  is still linear while that of  $\text{Bi}_2\text{Te}_3$  bends downwards, mainly due to the increase in the thermal conductivity due to bipolar effect.  $\text{Bi}_2\text{Te}_3$  reached a  $ZT$  of 1.04, while  $\text{Sb}_2\text{Te}_3$  reached to 1.20 at 440 K. The ternary phases where Bi and Sb are co-doped also showed favorable  $ZT$  values by increasing temperature, reaching about 0.75–0.90 at 523 K. The highest  $ZT$  of 1.37 is achieved for pure  $\text{Sb}_2\text{Te}_3$  sample at 523 K, likely due to its lower thermal conductivity that is suppressed by phonon point defect and phonon grain boundary scattering.

Fine thermoelectric powders as presented here are very valuable resources in the field of large area and flexible hybrid thermoelectric materials research. Many of the published works on hybrids have relied on commercial powders for making their films where micron size  $\text{Bi}_2\text{Te}_3$  material is either exfoliated via sonication, or ball-milled to reduce the particle size to prepare a good polymer nanoparticle hybrid. We attempted to show the ease of integration of as-made  $\text{Bi}_2\text{Te}_3$  material with PVDF and made few hybrid films using doctor blading method (see Figure S8). Our results show that stable dispersions can be easily made and films with reasonable transport properties can be obtained. A typical output from this film is n-type Seebeck ( $S = -46.12 \mu\text{V/K}$ ), and electrical conductivity of around 0.6013 S/m leading to a power factor of  $\text{PF} = 1.284 \times 10^{-3} \mu\text{W/mK}^2$  (see Figure S9). This value can be further improved by a careful interface engineering, the choice of the polymer, or changing the ratio of polymer vs. nanoparticle content in the hybrid films, which is one of the major directions of on-going research.

#### 4. Conclusions

Microwave assisted polyol synthesis was used to fabricate high-purity nanostructured  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  (where  $x$ : 0–2) thermoelectric materials in a matter of minutes, in a very highly reproducible manner.

Besides dramatically shortening the synthesis time, the quantities that can be fabricated per batch (>8 g) is unprecedented in other solution synthetic techniques. Fabricated materials exhibit hexagonal (or truncated hexagonal) platelet particle morphology in agreement with the rhombohedral crystal structure. Thermoelectric performance of  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  pellets were studied, where the highest ZT values of 1.04 (at 440 K) for  $x: 0$  ( $\text{Bi}_2\text{Te}_3$ ) and 1.37 (at 523 K) for  $x: 2$  ( $\text{Sb}_2\text{Te}_3$ ) were obtained, respectively, as n- and p-type TE materials. Ternary compositions, where  $x: 0.2$  and  $1.5$ , reached ZT of 0.8–0.9 at 523 K, which are also promising values for thermoelectric conversion. The large density of grain boundaries has been effective in reducing the thermal conductivity of the materials, thus making it possible to achieve ZT values in the order of 1 or above for pure phases of  $\text{Bi}_2\text{Te}_3$  and  $\text{Sb}_2\text{Te}_3$ . Furthermore, due to their small size and favorable surface chemistry, the as-made nanoparticles could easily be used to fabricate hybrid thermoelectric films. The devised MW-assisted synthesis technique, which is energy and resource effective resulting in thermoelectric materials with promising performance, opens avenues for large scale production of heat-energy harvesting materials. In addition, complete thermoelectric devices can be realized by combining our n-type and p-type  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  materials. Furthermore, the presented method is truly scalable and can be tuned for one-pot synthesis of even quaternary compositions such as  $(\text{BiSb})_2(\text{TeSe})_3$ .

**Supplementary Materials:** The following are available online at <http://www.mdpi.com/2076-3417/10/3/1178/s1>, Table S1. Literature review on microwave (MW)-assisted synthesis of  $\text{Bi}_2\text{Te}_3$ . Details of the synthetic route (solvent, reactor), reaction temperature-time; morphology of resultant structures; the quantity typically obtained according to described procedure in the articles; and the reported transport data. Figure S1. (a) Screenshot of the reactor for experimental design using multivessel rotor. (b) Multivessel high pressure rotor used for the MW-assisted synthesis of  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  nanoparticles. The system can take up to 15 reactors, which can be filled up to about 100 mL each, making it possible to synthesize high and reproducible quality nanostructures with a high yield, at a pilot scale, without needing any complicated or reflux systems. Figure S2. (a) Typical temperature (T)–time (t)–power (P) profile of the MW-assisted synthesis process for  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  samples. (b) As-synthesized sample transferred to a centrifuge vial and phase separation is easily achieved, yielding typically 2–3 g sample per reactor vessel. Four positions/vials (out of 15) during the MW-assisted reaction is sufficient to prepare 8–10 g of these materials within a reaction time of 2 min. Figure S3. Reaction mechanism showing the molecular structure of organic components used for the polyol synthesis of  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  nanoparticles. While  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  nanoparticles are formed, ethyleneglycol (EG) is converted to diacetyl (rxn (1)) and trioctylphosphine (TOP) is converted to trioctylphosphine oxide (TOPO) (rxn (2)). Figure S4. Estimated change of the lattice parameter  $a$  (or  $b$ ) by substitution of Sb ( $x$ ) in  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  system. Figure S5. (a) TEM micrograph of a typical as-made  $\text{Bi}_2\text{Te}_3$  nanoparticle with a lateral cross-sectional size of about 500 nm; (b) high resolution micrograph showing lattice fringes; and (c) Fourier-transform of the selected square region in (b) showing crystalline nature of  $\text{Bi}_2\text{Te}_3$  nanoplatelet. Figure S6. High resolution XPS Spectra of as-made  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$  samples, for the regions Bi 4f, Sb 3d, and Te 3d. Table S2. Transport properties for nanostructured bulk  $\text{Bi}_{2-x}\text{Sb}_x\text{Te}_3$ : mass density ( $\hat{\rho}$ ) in percentage of the theoretical density, total thermal conductivity ( $\kappa_{tot}$ ), Lorentz number ( $L_o$ ), electronic thermal conductivity ( $\kappa_{el}$ ), and lattice thermal conductivity ( $\kappa_{lat}$ ). Figure S7. Hybrid thermoelectric film ( $1.5 \times 5 \text{ cm}^2$ ) with Ag paste contact. The films are made by blending the as-made thermoelectric nanoparticles with the polymer PVDF. Figure S8. Microstructure of the as-made  $\text{Bi}_2\text{Te}_3$  nanoparticles–PVDF hybrid thermoelectric film at different magnifications. PVDF is visible as a dark shade in the rightmost micrograph, showing perfusing nanoparticles. Figure S9. Device characteristic of the large area film made with  $\text{Bi}_2\text{Te}_3$  nanoplatelets and PVDF. The resultant transport qualities for the film are  $S = -46.12 \mu\text{V/K}$ ,  $\sigma = 0.6013 \text{ S/m}$ , leading to a power factor of  $\text{PF} = 1.284 \times 10^{-3} \mu\text{W/mK}^2$ .

**Author Contributions:** M.S.T. conceived the presented idea. B.H., V.R., and M.S.T. contributed to conceptualization and methodology, material synthesis, and original draft preparation. B.H., H.B., and M.O. contributed to physicochemical characterization. B.H. and M.O. processed the material. S.B. and A.Y. validated the thermoelectric properties. M.K. and R.S. performed the XPS analysis. All authors discussed the results and reviewed and edited the manuscript. M.S.T. and M.J. supervised the work. All authors have read and agreed to the published version of the manuscript.

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## References

1. Boyer, A.; Cissé, E. Properties of thin film thermoelectric materials: Application to sensors using the Seebeck effect. *Mater. Sci. Eng. B* **1992**, *13*, 103–111.
2. Dresselhaus, M.S.; Dresselhaus, G.; Sun, X.; Zhang, Z.; Cronin, S.B.; Koga, T. Low dimensional thermoelectric materials. *Phys. Solid State* **1999**, *41*, 679–682. [[CrossRef](#)]
3. Hicks, L.D.; Dresselhaus, M.S. Thermoelectric figure of merit of a one-dimensional conductor. *Phys. Rev. B-Condens. Matter Mater. Phys.* **1993**, *47*, 16631–16634.
4. Hicks, L.D.; Dresselhaus, S.M. Effect of quantum-well structures on the thermoelectric figure of merit. *Phys. Rev.* **1993**, *47*, 12727–12731.
5. Venkatasubramanian, R.; Siivola, E.; Colpitts, T.; Quinn, B.O. Thin-film thermoelectric devices with high room-temperature figures of merit. *Nature* **2001**, *413*, 597–602. [[CrossRef](#)] [[PubMed](#)]
6. Lin, Y.M.; Sun, X.; Dresselhaus, M. Theoretical investigation of thermoelectric transport properties of cylindrical Bi nanowires. *Phys. Rev. B-Condens. Matter Mater. Phys.* **2000**, *62*, 4610–4623.
7. Mamur, H.; Bhuiyan, M.R.A.; Korkmaz, F.; Nil, M. A review on bismuth telluride (Bi<sub>2</sub>Te<sub>3</sub>) nanostructure for thermoelectric applications. *Renew. Sustain. Energy Rev.* **2018**, *82*, 4159–4169. [[CrossRef](#)]
8. Madavali, B.; Kim, H.S.; Lee, K.H.; Isoda, Y.; Gascoin, F.; Hong, S.J. Large scale production of high efficient and robust p-type Bi–Sb–Te based thermoelectric materials by powder metallurgy. *Mater. Des.* **2016**, *112*, 485–494. [[CrossRef](#)]
9. Serrano-Sánchez, F.; Gharsallah, M.; Nemes, N.M.; Biskup, N.; Varela, M.; Martínez, J.L.; Fernández-Díaz, M.T.; Alonso, J.A. Enhanced figure of merit in nanostructured (Bi,Sb)<sub>2</sub>Te<sub>3</sub> with optimized composition, prepared by a straightforward arc-melting procedure. *Sci. Rep.* **2017**, *7*, 1–10. [[CrossRef](#)]
10. Gharsallah, M.; Serrano-Sánchez, F.; Bermúdez, J.; Nemes, N.M.; Martínez, J.L.; Elhalouani, F.; Alonso, J.A. Nanostructured Bi<sub>2</sub>Te<sub>3</sub> Prepared by a Straightforward Arc-Melting Method. *Nanoscale Res. Lett.* **2016**, *11*, 1–7. [[CrossRef](#)]
11. Luo, Y.; Yang, J.; Jiang, Q.; Fu, L.; Xiao, Y.; Li, W.; Zhang, D.; Zhou, Z.; Cheng, Y. Melting and solidification of bismuth antimony telluride under a high magnetic field: A new route to high thermoelectric performance. *Nano Energy* **2015**, *15*, 709–718. [[CrossRef](#)]
12. Cao, Y.Q.; Zhao, X.B.; Zhu, T.J.; Zhang, X.B.; Tu, J.P. Syntheses and thermoelectric properties of Bi<sub>2</sub>Te<sub>3</sub> Sb<sub>2</sub>Te<sub>3</sub> bulk nanocomposites with laminated nanostructure. *Appl. Phys. Lett.* **2008**, *92*, 90–93.
13. Cao, Y.Q.; Zhu, T.J.; Zhao, X.B.; Zhang, X.B.; Tu, J.P. Nanostructuring and improved performance of ternary Bi–Sb–Te thermoelectric materials. *Appl. Phys. A Mater. Sci. Process.* **2008**, *92*, 321–324.
14. Jiang, Y.; Zhu, Y.J. Bi<sub>2</sub>Te<sub>3</sub> nanostructures prepared by microwave heating. *J. Cryst. Growth* **2007**, *306*, 351–355.
15. Shi, W.; Zhou, L.; Song, S.; Yang, J.; Zhang, H. Hydrothermal synthesis and thermoelectric transport properties of impurity-free antimony telluride hexagonal nanoplates. *Adv. Mater.* **2008**, *20*, 1892–1897.
16. Li, S.; Zhang, S.; He, Z.; Toprak, M.; Stiewe, C.; Muhammed, M.; Müller, E. Novel solution route synthesis of low thermal conductivity nanocrystalline bismuth telluride. *J. Nanosci. Nanotechnol.* **2010**, *10*, 7658–7662. [[CrossRef](#)]
17. Li, S.; Toprak, M.S.; Soliman, H.M.; Zhou, J.; Muhammed, M.; Platzek, D.; Müller, E. Fabrication of nanostructured thermoelectric bismuth telluride thick films by electrochemical deposition. *Chem. Mater.* **2006**, *18*, 3627–3633.
18. Yu, H.; Gibbons, P.C.; Buhro, W.E. Bismuth, tellurium, and bismuth telluride nanowires. *J. Mater. Chem.* **2004**, *14*, 595–602.
19. Li, S.; Liang, Y.; Qin, J.; Toprak, M.; Muhammed, M. Template electrodeposition of ordered bismuth telluride nanowire arrays. *J. Nanosci. Nanotechnol.* **2009**, *9*, 1543–1547.
20. Prieto, A.L.; Sander, M.S.; Martín-González, M.S.; Gronsky, R.; Sands, T.; Stacy, A.M. Electrodeposition of ordered Bi<sub>2</sub>Te<sub>3</sub> nanowire arrays. *J. Am. Chem. Soc.* **2001**, *123*, 7160–7161.
21. Xiao, F.; Yoo, B.; Kyu, H.L.; Myung, N.V. Synthesis of Bi<sub>2</sub>Te<sub>3</sub> nanotubes by galvanic displacement. *J. Am. Chem. Soc.* **2007**, *129*, 10068–10069.
22. Zhao, X.B.; Ji, X.H.; Zhang, Y.H.; Zhu, T.J.; Tu, J.P.; Zhang, X.B. Bismuth telluride nanotubes and the effects on the thermoelectric properties of nanotube-containing nanocomposites. *Appl. Phys. Lett.* **2005**, *86*, 062111.
23. Zhang, C.; Peng, Z.; Li, Z.; Yu, L.; Khor, K.A.; Xiong, Q. Controlled growth of bismuth antimony telluride Bi<sub>x</sub>Sb<sub>2-x</sub>Te<sub>3</sub> nanoplatelets and their bulk thermoelectric nanocomposites. *Nano Energy* **2015**, *15*, 688–696.

24. Harpeness, R.; Gedanken, A. Microwave-assisted synthesis of nanosized  $\text{Bi}_2\text{Se}_3$ . *New J. Chem.* **2003**, *27*, 1191–1193. [[CrossRef](#)]
25. Zhou, B.; Zhao, Y.; Pu, L.; Zhu, J. Microwave-assisted synthesis of nanocrystalline  $\text{Bi}_2\text{Te}_3$ . *Mater. Chem. Phys.* **2006**, *96*, 192–196.
26. Yao, Q.; Zhu, Y.; Chen, L.; Sun, Z.; Chen, X. Microwave-assisted synthesis and characterization of  $\text{Bi}_2\text{Te}_3$  nanosheets and nanotubes. *J. Alloys Compd.* **2009**, *481*, 91–95.
27. Pradhan, S.; Das, R.; Bhar, R.; Bandyopadhyay, R.; Pramanik, P. A simple fast microwave-assisted synthesis of thermoelectric bismuth telluride nanoparticles from homogeneous reaction-mixture. *J. Nanoparticle Res.* **2017**, *69*, 19.
28. Semenkovich, S.A.; Melekh, B.T. Thermodynamic Properties of  $\text{Bi}_2\text{Te}_3$ ,  $\text{Bi}_2\text{Se}_3$ ,  $\text{Sb}_2\text{Te}_3$  and  $\text{Sb}_2\text{Se}_3$ . In *Chemical Bonds in Solids*; Springer: New York, NY, USA, 1972; pp. 159–162.
29. Kosmulski, M. Isoelectric points and points of zero charge of metal (hydr)oxides: 50 years after Parks' review. *Adv. Colloid Interface Sci.* **2016**, *238*, 1–61. [[CrossRef](#)]
30. Crist, B.V. *Handbook of Monochromatic XPS Spectra—The Elements and Native Oxides*; Volume 1 in PDF; XPS International LLC: Mountain View, CA, USA, 1999.
31. Scheele, M.; Oeschler, N.; Veremchuk, I.; Reinsberg, K.G.; Kreuziger, A.M.; Kornowski, A.; Broekaert, J.; Klinke, C.; Weller, H. ZT enhancement in solution-grown  $\text{Sb}_{(2-x)}\text{Bi}_x\text{Te}_3$  nanoplatelets. *ACS Nano* **2010**, *4*, 4283–4291. [[CrossRef](#)]
32. Zhao, Y.; Dyck, J.S.; Hernandez, B.M.; Burda, C. Enhancing thermoelectric performance of ternary nanocrystals through adjusting carrier concentration. *J. Am. Chem. Soc.* **2010**, *132*, 4982–4983. [[CrossRef](#)]
33. Poudel, B.; Hao, Q.; Ma, Y.; Lan, Y.; Minnich, A.; Yu, B.; Yan, X.; Wang, D.; Muto, A.; Vashaee, D.; et al. High-thermoelectric performance of nanostructured bismuth antimony telluride bulk alloys. *Science* **2008**, *320*, 634–638. [[CrossRef](#)] [[PubMed](#)]
34. Xu, Z.; Wu, H.; Zhu, T.; Fu, C.; Liu, X.; Hu, L.; He, J.; He, J.; Zhao, X. Attaining high mid-Temperature performance in  $(\text{Bi,Sb})_2\text{Te}_3$  thermoelectric materials via synergistic optimization. *NPG Asia Mater.* **2016**, *8*, e302. [[CrossRef](#)]
35. Sehr, R.; Testardi, L.R. The optical properties of p-type  $\text{Bi}_2\text{Te}_3$ - $\text{Sb}_2\text{Te}_3$  alloys between 2–15 microns. *J. Phys. Chem. Solids* **1962**, *23*, 1219–1224. [[CrossRef](#)]
36. Hu, L.P.; Zhu, T.J.; Wang, Y.G.; Xie, H.H.; Xu, Z.J.; Zhao, X.B. Shifting up the optimum figure of merit of p-type bismuth telluride-based thermoelectric materials for power generation by suppressing intrinsic conduction. *NPG Asia Mater.* **2014**, *6*, e88. [[CrossRef](#)]
37. Charoenphakdee, A.; Yamanaka, S.; Snyder, G.J. Enhancement of Thermoelectric of the Electronic Density of States. *Science* **2008**, *321*, 1457–1461.
38. Wang, H.; Gibbs, Z.M.; Takagiwa, Y.; Snyder, G.J. Tuning bands of PbSe for better thermoelectric efficiency. *Energy Environ. Sci.* **2014**, *7*, 804–811. [[CrossRef](#)]

