Approximate Controllability of Sub-Diffusion Equation with Impulsive Condition

Lakshman Mahto 1, Syed Abbas 2,*, Mokhtar Hafayed 3 and Hari M. Srivastava 4,5,

1 Department of Science and Humanities, Indian Institute of Information Technology Dharwad, Hubli 580029, India; lm.iitmandi@gmail.com
2 School of Basic Sciences, Indian Institute of Technology Mandi, Mandi 175001, H.P., India
3 Laboratory of Applied Mathematics, Biskra University, Biskra 07000, Algeria; hafa.mokh@yahoo.com
4 Department of Mathematics and Statistics, University of Victoria, Victoria, BC V8W 3R4, Canada; harimsri@math.uvic.ca
5 Department of Medical Research, China Medical University Hospital, China Medical University,
Taichung 40402, Taiwan

* Correspondence: sabbas.iitk@gmail.com

Received: 7 January 2019; Accepted: 14 February 2019; Published: 17 February 2019

Abstract: In this work, we study an impulsive sub-diffusion equation as a fractional diffusion equation of order $\alpha \in (0, 1)$. Existence, uniqueness and regularity of solution of the problem is established via eigenfunction expansion. Moreover, we establish the approximate controllability of the problem by applying a unique continuation property via internal control which acts on a sub-domain.

Keywords: fractional diffusion equation; controllability; impulsive system; unique continuation property

MSC: 93B05; 34A08; 47H10; 49J15; 26A33; 45K05

1. Introduction

1.1. Fractional Diffusion Equations

A fractional diffusion equation of order $\alpha \in (0, 1)$ is obtained by rewriting a normal diffusion equation in integral form as

$$u(x,t) + \int_0^t Au(x,t)dt = u_0 + \int_0^t f(x,t)dt, \quad (x,t) \in \Omega \times (0, T). \quad (1)$$

Then, replacing the first of right-hand side (RHS) integral of Equation (1) by a Riemann-Liouville fractional integral, $I^\alpha$ of order $0 < \alpha < 1$, we get

$$u(x,t) + \int_0^t \frac{(t-s)^{\alpha-1}}{\Gamma(\alpha)} Au(x,t)dt = u_0 + \int_0^t f(x,t)dt, \quad (x,t) \in \Omega \times (0, T).$$

Now, differentiating the above equation on both sides with respect to $t$, we get the following fractional diffusion equation:

$$\partial_t u + \partial_t^{\alpha-a} u = f(x,t) \quad \text{in} \quad \Omega \times (0, T),$$

$$u = 0 \quad \text{on} \quad (\Gamma = \partial \Omega) \times (0, T),$$

$$u(\cdot, 0) = u_0 \quad \text{in} \quad \Omega. \quad (2)$$
If \( \alpha = 1 \), then Equation (2) is a classical diffusion equation. Equation (2) with \( 0 < \alpha < 1 \) is called the fractional diffusion equation. These equations appear in the model of anomalous diffusion in heterogeneous media. Anomalous diffusion is one of the most ubiquitous phenomena in nature; it has been observed in various fields of physical sciences, for example, surface growth, transport of fluid in porous media, two-dimensional rotating flow and diffusion of plasma. Because of such anomalies, the classical diffusion models can not be used to study the dynamics of such systems. In this situation, fractional derivatives extend the help and play a crucial role in characterizing such diffusion. The model corresponding to such derivative is called a fractional partial differential equation. From the continuous time random walk (CTRW) model, Metzler and Klafter [1] derived Equation (3) with \( 0 < \alpha < 1 \) as a macroscopic model.

### 1.2. Impulsive Partial Differential Equations

Impulsive partial differential equations are a very important class of differential equations. These equations arise from the modelling of various real world processes having memory and are subject to short time fluctuations. The theory of impulsive differential equation is very rich and wide. It is mainly due to the fact that the it inherit intrinsic difficulties of the problems. These kinds of equations have lots of applications in different branches of Science and Engineering. These kinds of equations arise naturally from several physical and natural processes like earthquakes and pulse vaccination strategy. For more information, we refer to [2–4] and references therein. For more theoretical work, one can see the interesting book by Bainov and Simeonov [5]. The authors Shun et al. in [6] consider second-order impulsive Hamiltonian systems and established the existence of infinitely many solutions.

### 1.3. Controllability

In mathematical control theory, controllability and optimal control are two important concepts. In controllability, one studies the steering of a dynamical system from a given initial state to any other state or in the neighborhood of the state under some admissible control input. The cases where target states are defined in a given subregion are particularly very important; this situation arises in many real world applications. The last few decades have seen tremendous work in the controllability problems for integer order systems. Several techniques have been developed for solving such problems [7,8]. It has been seen that mostly authors worked on the problems with hard constraints on the state or control. This is mainly due to its applicability and importance in various applications in optimal control. Moreover, many authors have studied controllability of the semilinear, partial evolution equations, we refer to [9–17] and references therein. In a very interesting paper [14], Kenichi Fujishiro and Masahiro Yamamoto consider a partial differential equations with fractional order time derivatives and established approximate controllability by interior control.

### 1.4. The Problem under Consideration

Let \( \Omega \) be a bounded domain of \( \mathbb{R}^d \) with \( C^2 \) boundary \( \Gamma = \partial \Omega \). We consider the following initial value/boundary value problem of an impulsive sub-diffusion equation of order \( \alpha \in (0, 1) \):

\[
\begin{aligned}
\partial_t u + \partial_i^{\alpha} A u &= f(x, t) \quad \text{in} \quad \Omega \times (0, T), \\
\Delta u(\cdot, t_i) &= I_i(u(\cdot, t_i)), \quad i = 1, 2, 3, \cdots, P, \\
u &= 0 \quad \text{on} \quad \Gamma \times (0, T), \\
u(\cdot, 0) &= u_0 \quad \text{in} \quad \Omega.
\end{aligned}
\]
In Equation (3), \( u = u(x,t) \) is the state to be controlled and \( f = f(x,t) \) is the control which is localized in a subdomain \( \omega \) of \( \Omega \). We will act by \( f \) to drive the initial state \( u_0 = u_0(x) \) to some target function \( u_1 = u_1(x) \). The operator \( A \) is a symmetric and uniformly elliptic operator. The details will be specified later; \( T > 0 \) is also a constant. Several problems in applications can be modeled by the above equation. Some of them are: thermal diffusion in media with fractional geometry, underground environmental problems, highly heterogeneous aquifer, etc. [18]. In this paper, we study approximate controllability for fractional partial differential equations with impulses. We say that Equation (3) is approximately controllable if, for any \( u_1 \in L^2(\Omega) \) and \( \epsilon > 0 \), there exists a control \( f \) such that the solution \( u \) of (3) satisfies
\[
\|u(\cdot, T) - u_1\|_{L^2(\Omega)} \leq \epsilon. \tag{4}
\]

This paper is divided into four sections. In Section 2, we study requisite function spaces and some important basic results. In Section 3, we analyse the mild solutions of the Equation (3) by eigenfunction expansion. Section 4 is devoted to the study of a dual system of (3) and to establish a unique continuation property. In the last section, we establish the proof of approximate controllability.

2. Preliminaries

In this section, we state a few function spaces, notations and results in order to establish our main results. For the smooth reading of the manuscript, we first define the following class of spaces (for more details, we refer to Adams [19], Mahto [12]):

\[
L^p[a,b] = \left\{ f: [a,b] \rightarrow \mathbb{R} \mid f \text{ is Lebesgue measurable and } \int_a^b |f(t)|^p dt < \infty \right\},
\]

\[
AC[a,b] = \left\{ f: [0,T] \rightarrow \mathbb{R} \mid f \text{ is absolutely continuous on } [a,b] \right\},
\]

\[
C[a,b] = \left\{ f: [0,T] \rightarrow \mathbb{R} \mid f \text{ is continuous on } [a,b] \right\},
\]

\[
L^p(\Omega) = \left\{ f: \Omega \rightarrow \mathbb{R} \mid f \text{ is Lebesgue measurable and } \int_{\Omega} |f(x)|^p dx < \infty \right\},
\]

\[
H^1(\Omega) = \left\{ f, \frac{\partial f}{\partial x_1}, \cdots, \frac{\partial f}{\partial x_d} \in L^2(\Omega) \right\},
\]

\[
H^1_0(\Omega) = \left\{ f : H^1(\Omega) \text{ and } f = 0 \text{ on } \Gamma \right\},
\]

\[
AC(0,T; L^2(\Omega)) = \left\{ f: [0,T] \rightarrow L^2(\Omega) \mid f \in AC([t_0,t_1]; L^2(\Omega)) \cup AC((t_i,t_{i+1}); L^2(\Omega)), i = 1, \right. \]
\[
2, \cdots, P, x(t) \in t_i^+ \text{ and } x(t) = x(t_{i-}^-)
\]

\[
PC(0,T; L^2(\Omega)) = \left\{ f: [0,T] \rightarrow L^2(\Omega) \mid f \in C([t_0,t_1]; L^2(\Omega)) \cup C((t_i,t_{i+1}); L^2(\Omega), i = 1,2, \right. \]
\[
\cdots, P, x(t) \in t_i^+ \text{ and } x(t) = x(t_{i-}^-)
\]

The functions and operators defined below are very standard in the fractional calculus. For more details, we refer to [20]:

1. Mittag-Leffler function by
\[
E_{\alpha,\beta}(z) := \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)}, \quad z \in \mathbb{C},
\]

where \( \alpha > 0 \) and \( \beta \in \mathbb{R} \) are arbitrary constants. We can directly verify that \( E_{\alpha,\beta}(z) \) is an entire function of \( z \in \mathbb{C} \). As for the Mittag–Leffler functions, we have the following lemma.
Lemma 1. Let $0 < \alpha < 2$ and $\beta \in \mathbb{R}$ be arbitrary and $\mu$ satisfy $\pi \alpha / 2 < \mu < \min\{\pi, \pi \alpha\}$. Then, there exists a constant $C = C(\alpha, \beta, \mu) > 0$ such that

$$|E_{\alpha, \beta}(z)| \leq \frac{C}{1 + |z|}, \quad \mu \leq |\arg(z)| \leq \pi.$$  \hfill (5)

2. Reimann-Liouville integrals: For $\alpha > 0$ and $f \in L^1(0, T)$, we define $\alpha$-th order forward and backward integrals of $f$ by

$$I_{0+}^\alpha f(t) := \frac{1}{\Gamma(\alpha)} \int_0^t (t - \tau)^{\alpha - 1} f(\tau) d\tau,$$

$$I_{T-}^\alpha f(t) := \frac{1}{\Gamma(\alpha)} \int_t^T (\tau - t)^{\alpha - 1} f(\tau) d\tau.$$  

In other words, the forward integral operators of $\alpha$-th order is the convolution with $t^{\alpha - 1}/\Gamma(\alpha)$ and consequently $I_{0+}^\alpha f$ also belongs to $L^1(0, T)$. The same argument is also valid for the backward integrals.

3. The Riemann-Liouville fractional derivatives: For $\alpha \in (0, 1)$, we define the forward and backward fractional derivatives of $f \in AC[0, T]$ by

$$\partial_+^\alpha f(t) := \frac{d}{dt} I_{0+}^{1-\alpha} h(t) = \frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_0^t (t - \tau)^{-\alpha} h(\tau) d\tau,$$

$$D_+^\alpha f(t) := \frac{1}{\Gamma(1-\alpha)} \left(-\frac{d}{dt}\right) \int_t^T (\tau - t)^{-\alpha} h(\tau) d\tau.$$  \hfill (6)

$$D_-^\alpha f(t) := \frac{1}{\Gamma(1-\alpha)} \left(-\frac{d}{dt}\right) \int_t^T (\tau - t)^{-\alpha} h(\tau) d\tau.$$  \hfill (7)

We also have the following lemmas for fractional integration by parts.

**Lemma 2.** Let $\alpha > 0$. If $f, g \in PC([0, T], L^2(\Omega))$, then

$$\int_0^T I_{0+}^\alpha f(t)g(t) dt = \int_0^T f(t) I_{T-}^\alpha g(t) dt.$$

**Proof.**

$$\int_0^T g(t) I_{0+}^\alpha f(t) dt = \int_0^T g(t) \int_0^t \frac{(t - s)^{\alpha - 1}}{\Gamma(\alpha)} f(s) ds dt$$

$$= \int_0^T f(t) \int_t^T \frac{(s - t)^{\alpha - 1}}{\Gamma(\alpha)} g(s) ds dt$$

(using Fubini theorem for change of order of integration.)

$$= \int_0^T f(t) I_{T-}^\alpha g(t) dt.$$

\qed

**Lemma 3.** Let $f \in PC(0, T), g \in C_0^\infty(0, T)$. Then, we have the following identity:

$$\int_0^T g(t) \partial_+^\alpha f(t) dt = \int_0^T f(t) D_+^\alpha g(t) dt.$$  \hfill (8)
Proof. By substituting the value of R-L fractional derivative, we obtain

\[
\begin{align*}
\int_0^T g(t) \frac{d}{dt} f(t) dt &= \int_0^T g(t) \frac{d}{dt} \int_0^t (t-s)^{-\alpha} f(s) ds dt \\
&= \left( \frac{g(t)}{\Gamma(1-\alpha)} \int_0^t (t-s)^{-\alpha} f(s) ds \right)_{t=0}^{t=T} - \int_0^T g'(t) \int_0^t (t-s)^{-\alpha} f(s) ds dt \\
&= \frac{g(T)}{\Gamma(1-\alpha)} \int_0^T (t-s)^{-\alpha} f(s) ds - \int_0^T g'(t) \int_0^t (t-s)^{-\alpha} f(s) ds dt \\
&= - \int_0^T g'(t) \int_0^t (t-s)^{-\alpha} f(s) ds dt \quad (\because g(T) = 0.) \\
&= - \int_0^T f(t) \frac{d}{dt} \int_t^T (s-t)^{1-\alpha} g'(s) ds dt \\
&= - \int_0^T f(t) \frac{d}{dt} \left( \frac{g(T)(T-t)^{-\alpha}}{\Gamma(1-\alpha)} - \frac{d}{dt} \int_t^T (s-t)^{1-\alpha} g'(s) ds \right) dt \\
&= \int_0^T f(t) \left( \frac{d}{dt} \left( \frac{g(s)(s-t)^{1-\alpha}}{\Gamma(2-\alpha)} \right) \right)_{s=t}^{s=T} - \frac{d}{dt} \int_t^T (s-t)^{1-\alpha} g'(s) ds dt \\
&= - \int_0^T f(t) \frac{d}{dt} \int_t^T (t-s)^{-\alpha} g(s) ds dt \\
&= \int_0^T f(t) D^\alpha_t g(t) dt.
\end{align*}
\]

\[\Box\]

3. Solution of Primal System

3.1. Representation of the Solution

To derive the representation, we first focus on \( t \in [0, t_1] \). We can rewrite (3) as

\[ \partial_t u + (\beta \ast Au) = f(\cdot, t), u(0) = u_0, \quad (9) \]

where \( \beta(t) = \frac{t^{\alpha-1}}{\Gamma(\alpha)} \) and \( Au = -\nabla^2 u \) is a symmetric, self-adjoint, uniformly elliptic operator with domain \( D(A) = H^2(\Omega) \cap H^1_0(\Omega) \), the spectrum of \( A \) is entirely composed of a countable number of eigenvalues and we can set with finite multiplicities:

\[ 0 < \lambda_1 \leq \lambda_2 \leq \cdots \leq \lambda_n \leq \cdots. \]
By \( \varphi_n \in H^2(\Omega) \cap H^1_0(\Omega) \), we denote the orthonormal eigenfunction corresponding to \( \lambda_n \):

\[
A \varphi_n = \lambda_n \varphi_n, \quad n = 1, 2, \ldots.
\]

Then, the sequence \( \{ \varphi_n \}_{n \in \mathbb{N}} \) is an orthonormal basis in \( L^2(\Omega) \). Since \( u(t) \in L^2(\Omega) \), we have

\[
u(t) = \sum_{j=1}^{\infty} u_j(t) \varphi_j,
\]

where \( u_j(t) = (u(t), \varphi_j) \) is the \( j \)th Fourier coefficient. Taking an inner product between (9) and \( \varphi_j \), we have an infinite number of linear integro-differential equations:

\[
\partial_t u_j(t) + \lambda_j (\beta \ast u_j) = f_j(\cdot, t), \quad (10)
\]

where \( f_j(\cdot, t) = (f(\cdot, t), \varphi_j) \) and \( u_{j0} = (u_0, \varphi_j) \).

Taking Laplace Transform both sides of (10), we get

\[
\mathcal{L}^{-1}\left(\frac{1}{z + \lambda_j^{1-a}}\right) = E_a(-\lambda_j^{t^a}). \quad (11)
\]

Now, the representation for \( u_j \) of (10) is given by

\[
u_j = E_a(-\lambda_j^{t^a}) u_{j0} + \int_0^\infty E_a(-\lambda_j(t-s)^a) f_j(\cdot, s) ds. \quad (12)
\]

Thus, a formal solution of (9) is given by

\[
u(t) = \mathcal{E}(t) u_0 + \int_0^t \mathcal{E}(t-s) f(\cdot, s) ds, \quad (13)
\]

where

\[
\mathcal{E}(t) u_0 = \sum_{j=1}^{\infty} E_a(-\lambda_j^{t^a})(u_0, \varphi_j) \varphi_j, \quad (14)
\]

\[
u(\cdot, t) = \begin{cases} \mathcal{E}(t) u_0 + \int_0^t \mathcal{E}(t-s) f(\cdot, s) ds, & t \in [0, t_1], \\ \mathcal{E}(t) u_0 + \sum_{i=1}^{t_1} \mathcal{E}(t-t_i) I_i(u(\cdot, t_i)) + \int_{t_i}^t \mathcal{E}(t-s) f(\cdot, s) ds, & t \in (t_i, t_{i+1}], \quad i = 1, \ldots, P. \end{cases} \quad (15)
\]
3.2. Weak Formulation

Rewriting the (3) in unified form, we get

\[
\begin{aligned}
\partial_t u + \partial_t^{-\alpha} Au &= f(x, t) + \sum_{1 \leq i < p} I_i(u(\cdot, t_i)) \delta(t - t_i) \quad \text{in} \quad \Omega \times (0, T), \\
u &= 0 \quad \text{on} \quad \Gamma \times (0, T), \\
u(\cdot, 0) &= u_0 \quad \text{in} \quad \Omega.
\end{aligned}
\]  

(16)

A weak formulation of (16) is to find a \( u \in \text{PC}(0, T; H^1_0(\Omega)) \) such that

\[
(\partial_t u, v) + (\partial_t^{-\alpha} Au, v) = (f, v) + \sum_{i} (I_i(u(\cdot, t_i)) \delta(t - t_i), v), v \in H^1_0(\Omega).
\]  

(17)

Thus, we have a variational form of (16) as follows:

\[
(\partial_t u, v)dt + a(u, v) = l(v),
\]  

(18)

where,

\[
a(u, v) = (\partial_t^{-\alpha} Au, v) = \int_{\Omega} \partial_t^{-\alpha} \nabla u \cdot \nabla v dx,
\]

\[
l(v) = (f, v)dt + \sum_{i} (I_i(u(t_i)) \delta(t - t_i), v),
\]

with the following conditions:

1. \( a(\cdot, \cdot) \) is bounded or continuous i.e. \( |a(u, v)|_{H^1_0(\Omega)} \leq C_1 \|u\|_{H^1_0(\Omega)} \|v\|_{H^1_0(\Omega)} \)
2. \( a(\cdot, \cdot) \) is coercive i.e. \( A(u, u) \geq C_2 \|u\|_{H^1_0(\Omega)} \)
3. \( l \) is continuous.

**Definition 1.** A function \( u : [0, T] \rightarrow H^1_0(\Omega) \) is called a weak solution of (3) if:

1. \( u \in L^2(0, T; H^1_0(\Omega)) \cap \text{PC}(0, T; H^1_0(\Omega)) \) and \( \partial_t u \in L^2(0, T; H^{-1}(\Omega)) \cap \text{PC}(0, T; H^{-1}(\Omega)) \),
2. For every \( v \in H^1_0(\Omega) \), \( u \) satisfies (18),
3. \( u(0) = u_0 \).

Based on the above analysis, we can now formulate the following two theorems.

**Theorem 1.** For every \( f \in L^2(0, T; H^{-1}(\Omega)) \) and \( u_0 \in H^1_0(\Omega) \), there exists a unique weak solution \( u \in L^2(0, T; H^1_0(\Omega)) \cap \text{PC}(0, T; H^1_0(\Omega)) \) of (3).

**Proof.** Existence and uniqueness of weak solution is followed by the Lax-Milgram theorem. \( \square \)

**Theorem 2.** For every \( f \in L^2(0, T; H^{-1}(\Omega)) \) and \( u_0 \in H^1_0(\Omega) \), there exists a unique mild solution \( u \in L^2(0, T; H^1_0(\Omega)) \cap \text{PC}(0, T; H^1_0(\Omega)) \) of (3) and given by (15).

4. Dual System

In order to establish approximate controllability, we also need to consider the dual system for (3), a similar strategy for partial differential equations of integer order (see Section 8 in [21] or Chapters 2 and 3 in [22] for example). The dual system for (3), which runs backward in time, is given by;
\[ \begin{cases} -\partial_t v + D_1^{1-\alpha} Av = 0 & \text{in } \Omega \times (0, T), \\ \Delta v(\cdot, t_i) = I_i^*(v(\cdot, t_i)), & i = 1, 2, 3, \ldots, P, \\ v = 0 & \text{on } \Gamma \times (0, T), \\ v(\cdot, T) = v_0 & \text{in } \Omega. \end{cases} \] \tag{19}

4.1. Solution of Dual System

**Proposition 1.** Let \( v_0 \in L^2(\Omega) \). Then, there exists a unique solution of (19) and the solution is given by

\[ v(x, t) = \sum_{n=1}^{\infty} (T - t)^{\alpha-1} E_{\alpha, \alpha} (-\lambda_n (T - t)^{\alpha})(v_0, \phi_n) \phi_n(x) \\
+ \sum_{i < T - t} \sum_{n=1}^{\infty} (T - t - t_i)^{\alpha-1} E_{\alpha, \alpha} (-\lambda_n (T - t - t_i)^{\alpha})(I_i^*(v(\cdot, t_i)), \phi_n) \phi_n(x) \] \tag{20}

and has the following estimate:

\[ \|v(\cdot, t)\|_{L^2(\Omega)} \leq C \left( (T - t)^{\alpha-1} \|v_0\|_{L^2(\Omega)} + P \|I_n\|_{L^2(\Omega)} \left( \sum_{i < T - t} (T - t - t_i)^{2\alpha-2} \right)^{\frac{1}{2}} \right), \] \tag{21}

where \( \|I_n\|_{L^2(\Omega)} = \sup_{1 \leq i \leq P} \{ \|I_i\|_{L^2(\Omega)} \} \).

Moreover, the mapping \( v : [0, T] \to L^2(\Omega) \) is analytically extended to \( S_T := \{ z \in \mathbb{C}; \text{Re } z < T \} \).

**Proof.** Here, we establish existence and uniqueness of solution of (19) for \( v_0 = 0 \).

Multiplying (19) with \( \phi_n \) and setting \( v_n(t) = (v(\cdot, t), \phi_n) \), we get

\[ \partial_t v_n(t) + \lambda_n \partial_1^{1-\alpha} v_n(t) + \sum_{i < T - t} (I_i^*(v(t_i)), \phi_n) = 0. \] \tag{22}

Since

\[ |v_n(t)|^2 \leq \sum |v_n(t)|^2 = \|v(\cdot, t)\|^2_{L^2(\Omega)} \to 0 \text{ as } t \to T, \]

we have

\[ v_n(T) = 0. \] \tag{23}

From existence and uniqueness of the solution of the fractional differential equation (see [12]), we get

\[ v_n(t) = 0, \quad n = 1, 2, 3, \ldots. \]

As \( \{ \phi_n \} \) is a complete orthonormal system, we have

\[ v = 0 \text{ in } \Omega \times (0, T). \]

Thus, Equation (19) has a unique solution.

Now, we show the estimate (21).
By (20), we have
\[
\|v(\cdot, t)\|_{L^2(\Omega)}^2 \\
\leq \left\| \sum_{n=1}^{\infty} (v_0, \varphi_n) (T - t)^{a-1} E_{a, \alpha} (-\lambda_n (T - t)^a) \varphi_n \right\|_{L^2(\Omega)}^2 \\
+ \left\| \sum_{t_1 < t - t_i} \sum_{j=1}^{m_k} (I^*_t(v(\cdot, t)), \varphi_n) \varphi_n (x) \right\|_{L^2(\Gamma)}^2 \\
= \sum_{n=1}^{\infty} \left| (v_0, \varphi_n) (T - t)^{a-1} E_{a, \alpha} (-\mu_k (T - t)^a) \right|^2 \\
+ \sum_{n=1}^{\infty} \left| \sum_{t_1 < t - t_i} (I^*_t(v(\cdot, t)), \varphi_n) (T - t - t_i)^{a-1} E_{a, \alpha} (-\mu_k (T - t - t_i)^a) \right|^2 \\
= C^2 \left( \sum_{n=1}^{\infty} |(v_0, \varphi_n)|^2 \right) (T - t)^{2\alpha - 2} + C^2 \sum_{t_1 < t - t_i} (T - t - t_i)^{2\alpha - 2} \left( \sum_{n=1}^{\infty} |(I^*_t(v(\cdot, t)), \varphi_n)|^2 \right).
\]

Therefore,
\[
\|v(\cdot, t)\|_{L^2(\Omega)} \leq C \left( (T - t)^{a-1} \|v_0\|_{L^2(\Omega)} + P \|I^*_t\|_{L^2(\Omega)} \left( \sum_{t_1 < t - t_i} (T - t - t_i)^{2\alpha - 2} \right) \right)^{\frac{1}{2}}.
\]

Next, we show the analyticity of \(v(\cdot, t)\) in \(t \in S_T\).

We note that \(E_{a, \alpha} (-\lambda_n z)\) is an entire function (see [20] for example) and consequently each \((T - z)^{a-1} E_{a, \alpha} (-\lambda_n (T - z)^a)\) is analytic in \(z \in S_T\). Therefore, \(\sum_{n=1}^{\infty} (v_0, \varphi_n) (T - z)^{a-1} E_{a, \alpha} (-\lambda_n (T - z)^a) \varphi_n\) in \(S_T\).

If we fix \(\delta > 0\) arbitrarily, then, for \(z \in \mathbb{C}\) with \(\text{Re } z \leq T - \delta\), we have
\[
\left\| \sum_{n=M}^{N} (v_0, \varphi_n) (T - z)^{a-1} E_{a, \alpha} (-\lambda_n (T - z)^a) \varphi_n \right\|_{L^2(\Omega)}^2 \\
= \sum_{n=M}^{N} \left| (v_0, \varphi_n) (T - z)^{a-1} E_{a, \alpha} (-\lambda_n (T - z)^a) \right|^2 \\
\leq C \sum_{n=M}^{N} |(v_0, \varphi_n)|^2 |T - z|^{2\alpha - 2} \\
\leq C \delta^{2\alpha - 2} \sum_{n=M}^{N} |(v_0, \varphi_n)|^2 \to 0 \quad \text{as } M, N \to \infty.
\]

That is, (20) is uniformly convergent in \(\{ z \in \mathbb{C}; \text{Re } z \leq T - \delta \}\). Hence, \(v(\cdot, t)\) is also analytic in \(t \in S_T\). \( \square \)

4.2. Unique Continuation Property

**Proposition 2.** Let \(\omega\) be open in \(\Omega\) and \(v_0 \in L^2(\Omega)\). If a solution \(v \in PC(0, T; H^2(\Omega) \cap H^1_0(\Omega))\) be the solution of (19) vanishing in \(\omega \times (0, T)\), then \(v = 0\) in \(\Omega \times (0, T)\).
Proof. Since \( v(x,t) = 0 \) in \( \omega \times (0,T) \) and \( v : [0,T) \to L^2(\Gamma) \) can be analytically extended to \( S_T := \{ z \in \mathbb{C}; \text{Re} \, z < T \} \), we have

\[
v(x,t) = \sum_{n=1}^{\infty} (v_{\partial}(\varphi_n)(T-t)^{a-1}E_{a,a}(-\lambda_n(T-t)^a)\varphi_n(x) + \sum_{t<T} \sum_{n=1}^{\infty} (I^t_n(v(\cdot,t)),\varphi_n)(T-t-t_i)^{a-1}E_{a,a}(-\lambda_n(T-t-t_i)^a)\varphi_n(x) = 0, \quad x \in \omega, \ t \in (-\infty,T). \tag{24}\]

Let \( \{\mu_k\}_{k \in \mathbb{N}} \) be all spectra of \( L \) without multiplicities and we denote by \( \varphi_{\mu_j} \) an orthonormal basis of \( \text{Ker}(\mu_k-L) \). By using these notations, we can rewrite (24) by

\[
v(x,t) = \sum_{k=1}^{\infty} \sum_{j=1}^{m_k} (v_{\partial}(\varphi_{\mu_j})(x)(T-t)^{a-1}E_{a,a}(-\mu_k(T-t)^a) + \sum_{t<T} \sum_{j=1}^{m_k} (I^t_n(v(\cdot,t)),\varphi_{\mu_j})(x)(T-t-t_i)^{a-1}E_{a,a}(-\mu_k(T-t-t_i)^a) = 0, \quad x \in \omega, \ t \in (-\infty,T). \tag{25}\]

Then, for any \( z \in \mathbb{C} \) with \( \text{Re} \, z = \xi > 0 \) and \( N \in \mathbb{N} \), we have

\[
\left\| \sum_{k=1}^{N} \left( \sum_{j=1}^{m_k} (v_{\partial}(\varphi_{\mu_j})(x)) e^{i(t-T)}(T-t-t_i)^{a-1}E_{a,a}(-\mu_k(T-t)^a) \right) \right\|_{L^2(\Gamma)}^2 \\
= \sum_{k=1}^{N} \left( \sum_{j=1}^{m_k} (v_{\partial}(\varphi_{\mu_j})(x))^2 \right) e^{2\xi(t-T)}(T-t-t_i)^{a-1}E_{a,a}(-\mu_k(T-t)^a) \leq C^2 e^{2\xi(t-T)}(T-t-t_i)^{2a-2} \| v_\partial \|_{L^2(\Gamma)}
\]

and

\[
\left\| \sum_{k=1}^{N} \sum_{t<T-t_i} \sum_{j=1}^{m_k} (I^t_n(v(\cdot,t)),\varphi_{\mu_j})(x)(T-t-t_i)^{a-1}E_{a,a}(-\mu_k(T-t-t_i)^a) \right\|_{L^2(\Gamma)}^2 \\
= \sum_{k=1}^{N} \sum_{t<T-t_i} \sum_{j=1}^{m_k} (I^t_n(v(\cdot,t)),\varphi_{\mu_j})(x)^2 e^{2\xi(t-T)}(T-t-t_i)^{a-1}E_{a,a}(-\mu_k(T-t-t_i)^a) \leq C^2 e^{2\xi(t-T)}(T-t-t_i)^{2a-2} \| I^t_n \|_{L^2(\Gamma)} \sum_{t<T-t_i} (T-t-t_i)^{2a-2},
\]

where \( \| I^t_n \|_{L^2(\Gamma)} = \sup_{1 \leq i < p} \left\{ |t_i| \right\} \| t_i \|_{L^2(\Gamma)} \). Therefore,

\[
\left\| \sum_{k=1}^{N} \left( \sum_{j=1}^{m_k} (v_{\partial}(\varphi_{\mu_j})(x)) e^{i(t-T)}(T-t)^{a-1}E_{a,a}(-\mu_k(T-t)^a) \right) \right\|_{L^2(\Gamma)}^2 \\
\leq C e^{\xi(t-T)}(T-t)^{a-1} \| v_\partial \|_{L^2(\Gamma)}
\]
and
\[
\left\| \sum_{k=1}^{N} \sum_{t_{i} < t} \left( \sum_{j=1}^{m_k} (I^*_j(v(\cdot, t_i)), \varphi_{kj}(x)) \right) (T-t)_{i}^{a-1}E_{a,a}(-\mu_k(T-t)_{i}^{a}) \right\|_{L^2(\Gamma)} \leq PCe^{\delta(T-t)} \| I^n \|_{L^2(\Omega)} \left( \sum_{t>T-t_i} (T-t)_{i}^{2a-2} \right)^{\frac{1}{2}}.
\]

The right-hand sides of the two inequalities above are integrable on \((-\infty, T)\):
\[
\int_{-\infty}^{T} e^{\delta(T-t)}(T-t)^{a-1} dt = \frac{\Gamma(a)}{\delta^a}
\]
and
\[
\int_{-\infty}^{T-t_i} e^{\delta(T-t)}(T-t-t_i)^{a-1} dt = \int_{0}^{\infty} e^{\delta t} t^{a-1} dt = \frac{\Gamma(a)}{\delta^a}.
\]

Hence, the Lebesgue theorem yields that
\[
\int_{-\infty}^{T} e^{\delta(T-t)} \left( \sum_{k=1}^{\infty} \sum_{j=1}^{m_k} (v_0, \varphi_{kj}(x)) \right) (T-t)_{i}^{a-1}E_{a,a}(-\mu_k(T-t)_{i}^{a}) dt + \int_{-\infty}^{T-t_i} e^{\delta(T-t-T)} \times \left( \sum_{\lambda < T-t_i} \sum_{n=1}^{\infty} \sum_{j=1}^{m_k} (I^*_n(v(\cdot, t_i)), \varphi_{kj}(x)) (T-t-t_i)_{i}^{a-1}E_{a,a}(-\mu_k(T-t-t_i)_{i}^{a}) \right) dt
= \sum_{k=1}^{\infty} \sum_{j=1}^{m_k} \frac{(v_0 + \sum_{\lambda < T-t_i} I^*_n(v(\cdot, t_i)), \varphi_{kj}(x))}{z^a + \mu_k} \quad \text{a.e. } x \in \Omega, \text{ Re } z > 0,
\]
where we have used the Laplace transform formula;
\[
\int_{0}^{\infty} e^{-z t} t^{a-1} dt = \frac{1}{z^a + \mu_k}, \quad \text{Re } z > 0
\]
(see (1.80) in p. 21 of [20]). By (25) and (26), we have
\[
\sum_{k=1}^{\infty} \sum_{j=1}^{m_k} \frac{(v_0 + \sum_{\lambda < T-t_i} I^*_n(v(\cdot, t_i)), \varphi_{kj}(x))}{z^a + \mu_k} \varphi_{kj}(x) = 0, \quad \text{a.e. } x \in \omega, \text{ Re } z > 0,
\]
that is,
\[
\sum_{k=1}^{\infty} \sum_{j=1}^{m_k} \frac{(v_0 + \sum_{\lambda < T-t_i} I^*_n(v(\cdot, t_i)), \varphi_{kj}(x))}{\eta + \mu_k} \varphi_{kj}(x) = 0, \quad \text{a.e. } x \in \omega, \text{ Re } \eta > 0.
\]
By using analytic continuation in \(\eta\), we have
\[
\sum_{k=1}^{\infty} \sum_{j=1}^{m_k} \frac{(v_0 + \sum_{\lambda < T-t_i} I^*_n(v(\cdot, t_i)), \varphi_{kj}(x))}{\eta + \mu_k} \varphi_{kj}(x) = 0, \quad \text{a.e. } x \in \omega, \eta \in \mathbb{C} \setminus \{-\mu_k\}_{k \in \mathbb{N}}.
\]
(27)

Then, we can take a suitable disk which includes \(-\mu\ell\) and does not include \(-\mu_k\) for all \(k \neq \ell\).

By integrating (27) in the disk, we have
\[
\sum_{j=1}^{m_\ell} (v_0 + \sum_{t_i < T-t_i} I^*_j(v(\cdot, t_i)), \varphi_{\ell j}) \varphi_{\ell j}(x) = 0, \quad \text{a.e. } x \in \omega.
\]
By setting \(\tilde{\varphi}_\ell := \sum_{j=1}^{m_\ell} (v_0 + \sum_{t_i < T-t_i} I^*_j(v(\cdot, t_i)), \varphi_{\ell j}) \varphi_{\ell j}(x),\) we have
\[
(A - \mu_\ell)\tilde{\varphi}_\ell = 0 \quad \text{in } \Omega \quad \text{and} \quad \tilde{\varphi}_\ell = 0 \quad \text{on } \omega.
\]
Theorem 3. Let (see [23,24]) implies

If the conclusion of Theorem satisfies (3)

for (19) is equivalent to the approximate controllability for (3) stated in Theorem 3.

Therefore, the unique continuation result for eigenvalue problem of elliptic operator (see [23,24]) implies

\[ \tilde{\varphi}_\ell(x) = \sum_{j=1}^{m_\ell} (v_0 + \sum_{t<T-t_j} I^*_t(\varphi(\cdot,t_j)), \varphi_{\ell j}) \varphi_{\ell j}(x) = 0, \quad x \in \Omega \]

for each \( \ell \in \mathbb{N} \). Since \( \{ \varphi_{\ell j} \}_{1 \leq j \leq m_\ell} \) is linearly independent in \( \Omega \), we see that

\[ (v_0 + \sum_{t<T-t_j} I^*_t(\varphi(\cdot,t_j)), \varphi_{\ell j}) = 0, \quad 1 \leq j \leq m_\ell, \quad \ell \in \mathbb{N}. \]

This implies \( v = 0 \) in \( \Omega \times (0,T) \). \( \square \)

5. Approximate Controllability

In this section, we complete the proof of our main theorems.

**Theorem 3.** Let \( 0 < \alpha < 1 \) and \( \omega \) be an open set in \( \Omega \). Then, Equation (3) is approximately controllable for arbitrarily given \( T > 0 \). That is,

\[ \{ u(\cdot,T); f \in C_0^\infty(\omega \times (0,T)) \} = L^2(\Omega), \tag{28} \]

where \( u \) is the solution to (3) and the closure on the left-hand side is taken in \( L^2(\Omega) \).

We start the proof with a lemma.

**Lemma 4.** If the conclusion of Theorem (3) is true for \( u_0 = 0 \), then it is true for any \( u_0 \in H_0^1(\Omega) \).

**Proof.** Let \( u_0 \in H_0^1(\Omega) \) and \( u_T \in L^2(\Omega) \). Let \( \varepsilon > 0 \). Let us introduce \( \bar{u} \) the (mild) solution of

\[
\begin{aligned}
\dot{u}_t + \partial_1^{1-\alpha} A \bar{u} &= 0 \quad (x,t) \in \Omega \times (0,T), \\
\Delta u(\cdot,t_i) &= I_i(u(\cdot,t_i)), \quad i = 1,2,3,\ldots,P, \\
\bar{u}(x,t) &= 0, \quad t \in \Gamma \times (0,T), \\
\bar{u}(x,0) &= u_0(x), \quad x \in \Omega.
\end{aligned}
\]

Then, \( \bar{u}(T) \in L^2(\Omega) \). Therefore, using the assumption of Lemma 4, there exists \( f \in C_0^\infty(\omega \times (0,T)) \) such that the solution \( w \) of

\[
\begin{aligned}
\dot{w} + \partial_1^{1-\alpha} A w &= f(x,t) \quad (x,t) \in \Omega \times (0,T), \\
\Delta u(\cdot,t_i) &= I_i(w(\cdot,t_i)), \quad i = 1,2,3,\ldots,P, \\
w(1,t) &= 0, \quad t \in \Gamma \times (0,T), \\
w(x,0) &= 0, \quad x \in \Omega,
\end{aligned}
\]

satisfies

\[ \| w(T) - (u_T - \bar{u}(T)) \|_{L^2(\Omega)} \leq \varepsilon. \]

One can easily see that \( u(T) = w(T) + \bar{u}(T) \), so that the proof of Lemma 4 is achieved. \( \square \)

We now assume that \( u_0 \equiv 0 \).

In order to complete the proof of Theorem 3, we will see that the unique continuation property for (19) is equivalent to the approximate controllability for (3) stated in Theorem 3.
Proof. Let \( u \) be a solution of (3) for \( f \in C^0_0(\omega \times (0, T)) \) and let \( v \) be a solution of (19) for \( v_0 \in L^2(\Omega) \). Then, we see that

\[
0 = \int_0^T \int_\Omega (\partial_t u + \partial_1^{1-a} Au - f) v dx dt
\]

\[
= \int_0^T \int_\Omega (\partial_t u) v dx dt + \int_0^T \int_\Omega (\partial_1^{1-a} Au) v dx dt - \int_0^T \int_\Omega \sum_{1 \leq i \leq p} I_i(u(t_i)) \delta(t - t_i) v dx dt.
\]

In the above equation, the first term is calculated as follows:

\[
\int_0^{T-\delta} \int_\Omega (\partial_t u) v dx dt = \int_0^{T-\delta} \int_\Omega (\partial_t u) v dx dt
\]

\[
= \int_\Omega u v|_{t=0}^{t=T-\delta} dx - \int_0^{T-\delta} \int_\Omega u(\partial_t v) dx dt
\]

\[
= \int_\Omega u(\cdot, T-\delta) v_0 dx + \int_0^{T-\delta} \int_\Omega u(\partial_t v) dx dt.
\]

Here, we have used the integration in \( t \) by parts and the initial conditions in (3) and (19).

In terms of \( u \in PC(0, T; H^2(\Omega)) \) and \( v \in PC(0, T; H^2(\Omega) \cap H^1_0(\Omega)) \), we apply the Green formula to the second term, we have

\[
\int_0^{T-\delta} \int_\Omega (\partial_1^{1-a} Au) v dx dt = \int_0^{T-\delta} \int_\Omega (\partial_1^{1-a} u)(Au) v dx dt + \int_0^{T-\delta} \int_\Gamma \left( \frac{\partial v}{\partial \nu_A} - \frac{\partial u}{\partial \nu_A} v \right) d\sigma_A dt
\]

\[
= - \int_0^{T-\delta} \int_\Omega u(D_1^{1-a} Av) v dx dt.
\]

In the above calculation, we have used boundary conditions in (3) and (19). Therefore, we have

\[
0 = \int_0^{T-\delta} \int_\Omega (\partial_t u) v dx dt + \int_0^{T-\delta} \int_\Omega (\partial_1^{1-a} Au) v dx dt - \int_0^{T-\delta} \int_\Omega f v dx dt
\]

\[
= \int_\Omega u(\cdot, T-\delta) v_0 dx + \int_0^{T-\delta} \int_\Omega u(\partial_t v) dx dt - \int_0^{T-\delta} \int_\Omega u(D_1^{1-a} Av) v dx dt
\]

\[
- \int_0^{T-\delta} \int_\Omega f v dx dt - \int_0^{T-\delta} \int_\Omega \sum_{1 \leq i \leq p} I_i(u(t_i)) \delta(t - t_i) v dx dt
\]

\[
= \int_\Omega u(\cdot, T-\delta) v_0 dx + \int_0^{T-\delta} \int_\Omega u(\partial_t v - D_1^{1-a} Av) v dx dt
\]

\[
- \int_0^{T-\delta} \int_\Omega f v dx dt - \int_0^{T-\delta} \int_\Omega \sum_{1 \leq i \leq p} I_i(u(t_i)) v dx
\]

\[
= \int_\Omega u(\cdot, T-\delta) v_0 dx - \int_0^{T-\delta} \int_\Omega f v dx dt - \sum_{1 \leq i \leq p} \int_\Omega I_i(u(t_i)) v dx.
\]

Since \( u \in PC([0, T], L^2(\Omega)) \) and \( v(\cdot, T) = v_0 \) and taking \( \delta \to 0 \), we get

\[
\int_\Omega u(\cdot, T) v_0 dx - \sum_{1 \leq i \leq p} \int_\Omega I_i(u(t_i)) v dx = \int_0^T \int_\Omega f v dx dt.
\]
In order to prove density of \( \{u(\cdot, T); f \in C_0^\infty(\omega \times (0, T))\} \) in \( L^2(\Omega) \), we have to show that, if \( v_0 \in L^2(\Omega) \) satisfies
\[
(u(\cdot, T), v_0) = \int_{\Omega} u(\cdot, T)v_0dx = 0
\]
for any \( f \in C_0^\infty(\omega \times (0, T)) \), then \( v_0 \equiv 0 \). This can be shown as follows: we have
\[
\int_0^T \int_{\Omega} f(v_0)dxdt = 0
\]
for any \( f \in C_0^\infty(\omega \times (0, T)) \). Then, by the fundamental theorem of the calculus of variations, we have
\[
v(x, t) = 0, (x, t) \in \omega \times (0, T).
\]
By proposition (2), we have
\[
v(x, t) = 0, (x, t) \in \Omega \times (0, T).
\]
By uniqueness of the solution of (1),
\[
v_0(x) = 0, x \in \Omega,
\]
which gives \( \{u(\cdot, T); f \in C_0^\infty(\omega \times (0, T))\} \perp = \{0\} \). Hence, \( \{u(\cdot, T); f \in C_0^\infty(\omega \times (0, T))\} \) is dense in \( L^2(\Omega) \).

Thus, the proof of Theorem (3) is completed. \( \square \)

6. Example

Example 1. Consider the following relaxations’ oscillation equation with fractional order given by
\[
\frac{\partial}{\partial t} u(x, t) = \frac{\partial^{1-\alpha}}{\partial t^{1-\alpha}} u(x, t) + f(x, t), \quad t \in (0, T), \quad x \in \Omega = (0, \pi),
\]
\[
u(0, t) = u(\pi, t) = 0 \quad t \in (0, T),
\]
\[
u(x, 0) = u_0, \quad x \in (0, \pi),
\]
\[
\begin{align*}
\Delta u(x, t_k) & = -u(x, t_k) \quad k = 1, 2, \cdots, N.
\end{align*}
\]

Now, consider the corresponding system Let \( u(t)x = u(x, t) \) and assume \( f(x(t), t) \) to be a continuous function with respect to \( t \) that satisfies the Lipschitz condition in \( x \). Define the operator \( A = \frac{\partial^2}{\partial x^2} \) with domain
\[
D(A) = \{ x \in L^2(0, \pi) : x, x' \text{ are absolutely continuous and } x, x', x'' \in L^2(0, \pi) \}.
\]

It is well known that for \( \alpha = 1 \), sectorial operator, \( A = \frac{\partial^2}{\partial x^2} \) generates an analytic semigroup and for \( \alpha = 2 \), sectorial operator, \( A = \frac{\partial^2}{\partial x^2} \) generates a cosine family of operators.

Using the above notation, now consider the following system
\[
\begin{align*}
\frac{\partial u}{\partial t} & = \frac{\partial^{1-\alpha}}{\partial t^{1-\alpha}} Au, \quad t \in (0, T), \quad x \in \Omega = (0, \pi),
\end{align*}
\]
\[
\begin{align*}
u(0, t) & = u(\pi, t) = 0 \quad t \in (0, T),
\end{align*}
\]
\[
\begin{align*}
u(x, 0) & = 0, \quad x \in (0, \pi),
\end{align*}
\]
\[
\begin{align*}
\Delta u(x, t_k) & = -u(x, t_k) \quad k = 1, 2, \cdots, N.
\end{align*}
\]
The above problem can be posed as an abstract problem on $X = L^2(0, \pi) = U$, and hence it has a unique solution. Hence under the assumption of Theorem, the problem is approximately controllable.

**Example 2.** By choosing the function $\cos(t^2)\exp(-t)$, we get the following relaxation oscillation equation with fractional order given by

$$\partial_{1.8}^1 u(t) + Au(t) = \cos(t^2)\exp(-t), u(0) = 1, u'(0) = 1,$$

(33)

where $A$ is the operator mentioned above.

The graphical illustration of Example 2 is depicted in the Figure 1.

![Graphical Illustration](image)

**Figure 1.** Comparison of solution of (33) with varied relaxation coefficients, $A = 1, 2$ and 3.

**7. Discussion**

This paper presents a fractional sub-diffusion equation of an impulsive system (3) and its dual (19). The unique continuation Property 2 of the dual system plays a crucial role in the proof of our main result, approximate controllability Theorem 3 of the primal system with an interior control acts on a sub-domain. As an example, the approximate controllability of a fractional relaxation-oscillation equation is discussed and simulated for different relaxation coefficients.

**Author Contributions:** Conceptualization, L.M. and S.A.; methodology, S.A. and M.H.; software, M.H.; validation, S.A. and H.M.S.; formal analysis, L.M. and H.M.S.; writing—original draft preparation, S.A. and M.H.; writing—review and editing, S.A. and H.M.S.; supervision, S.A. and H.M.S.

**Funding:** This research received no external funding.

**Acknowledgments:** The authors are thankful to the anonymous reviewers for their careful reading of the manuscript and constructive comments and suggestions. Lakshman Mahto would like to thank The Institute of Mathematical Sciences, Chennai, for support and hospitality during the postdoctoral work, where this work was initiated.

**Conflicts of Interest:** The authors declare no conflict of interest.

**References**


